

Seismic tomography and the dilemma of the Earth's heat budget.



Guust Nolet
Géosciences Azur

The heat flux problem (I)



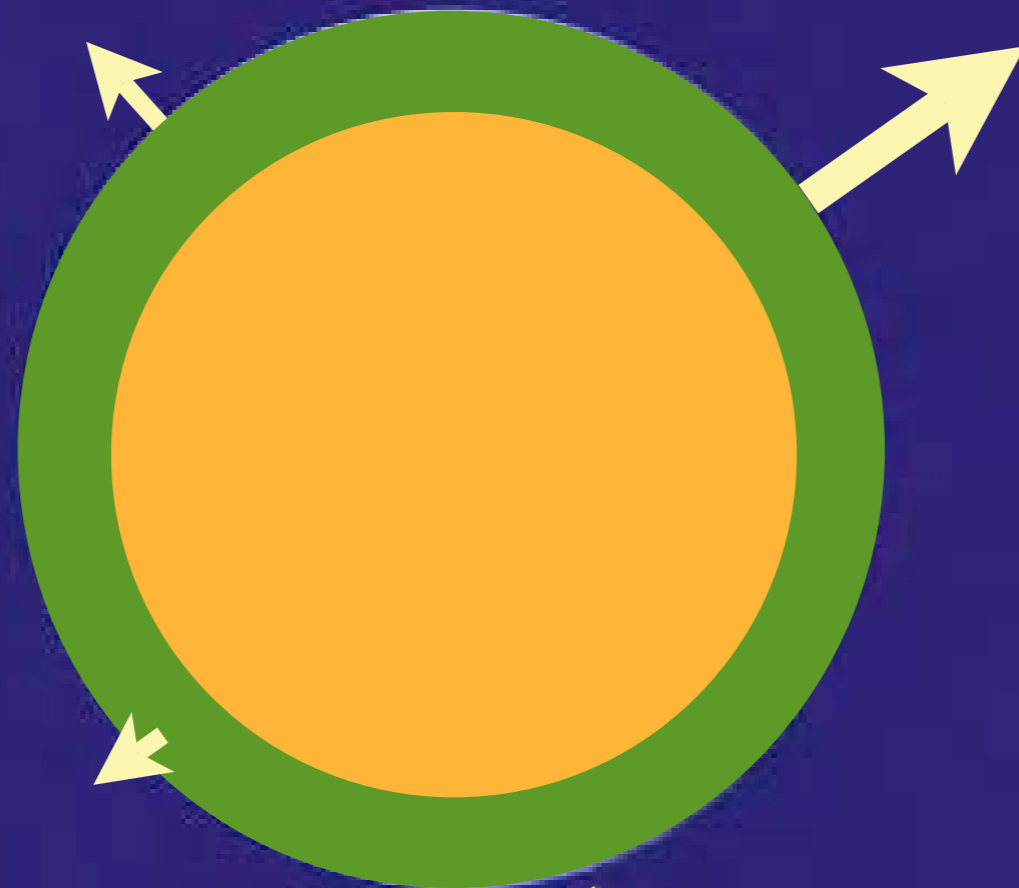
-8 TW from continents

44 TW



-8 TW from continents

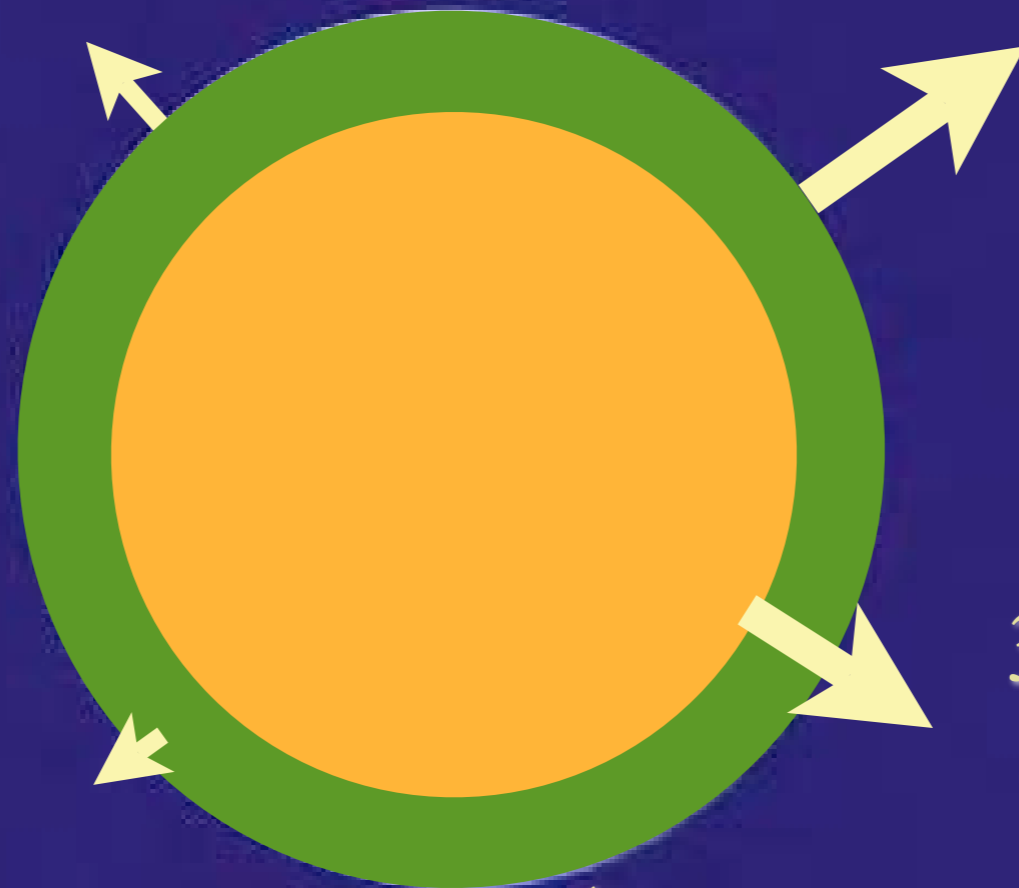
44 TW



-5 TW from upper mantle

-8 TW from continents

44 TW



-5 TW from upper mantle

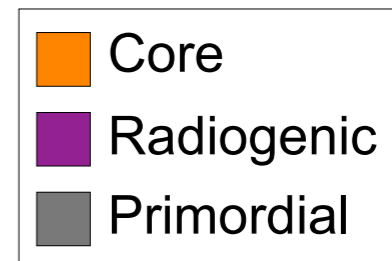
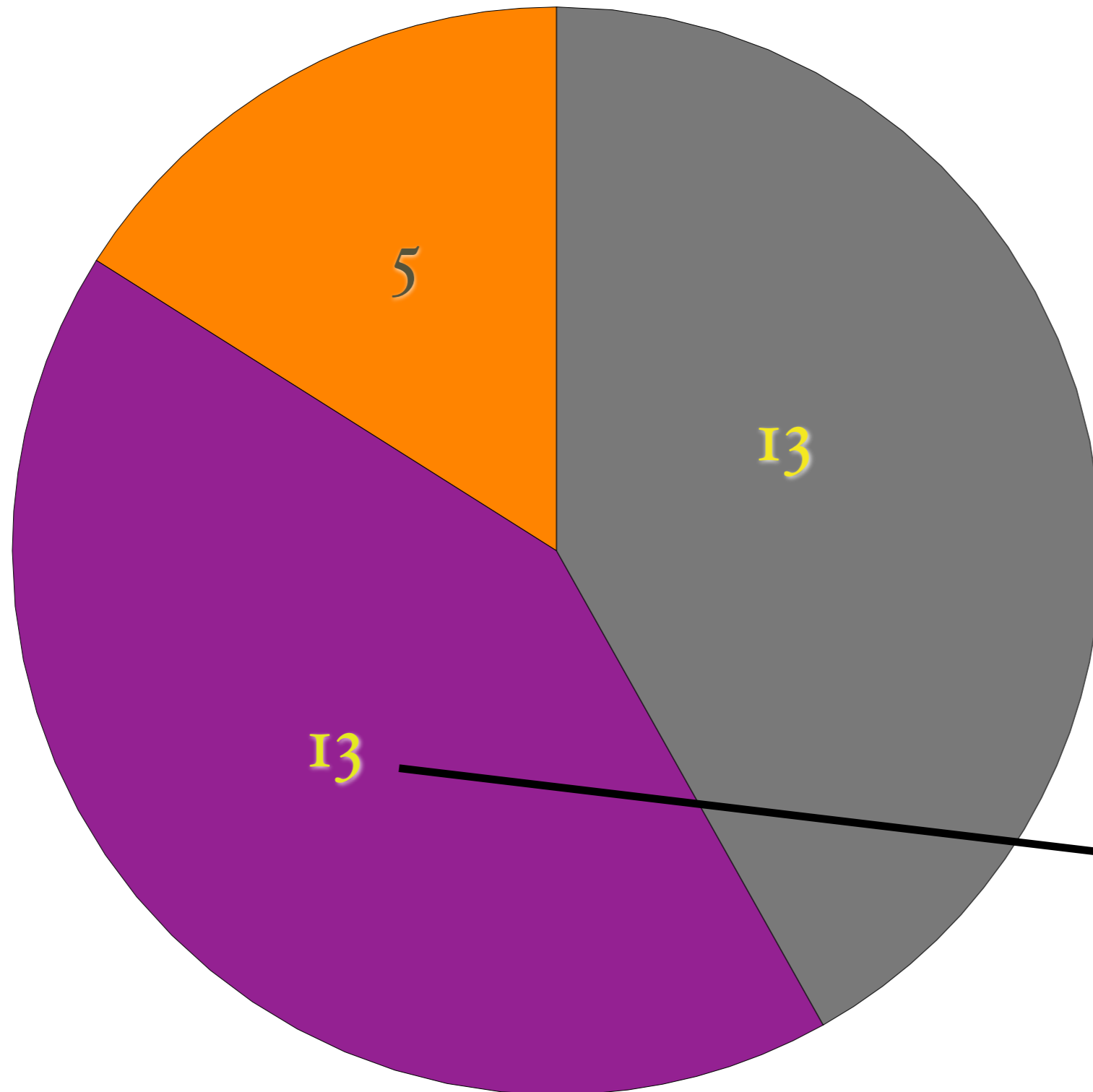
31 TW
across
660

Getting 31 TW across 660

- Whole-mantle convection
- Upward advection of hot rock (*plumes*)
- Downward advection of cold rock (*slabs*)
- Conduction (*thermal boundary layer or TBL*)

Question: Does a TBL exclude advection??

Mantle heat flux

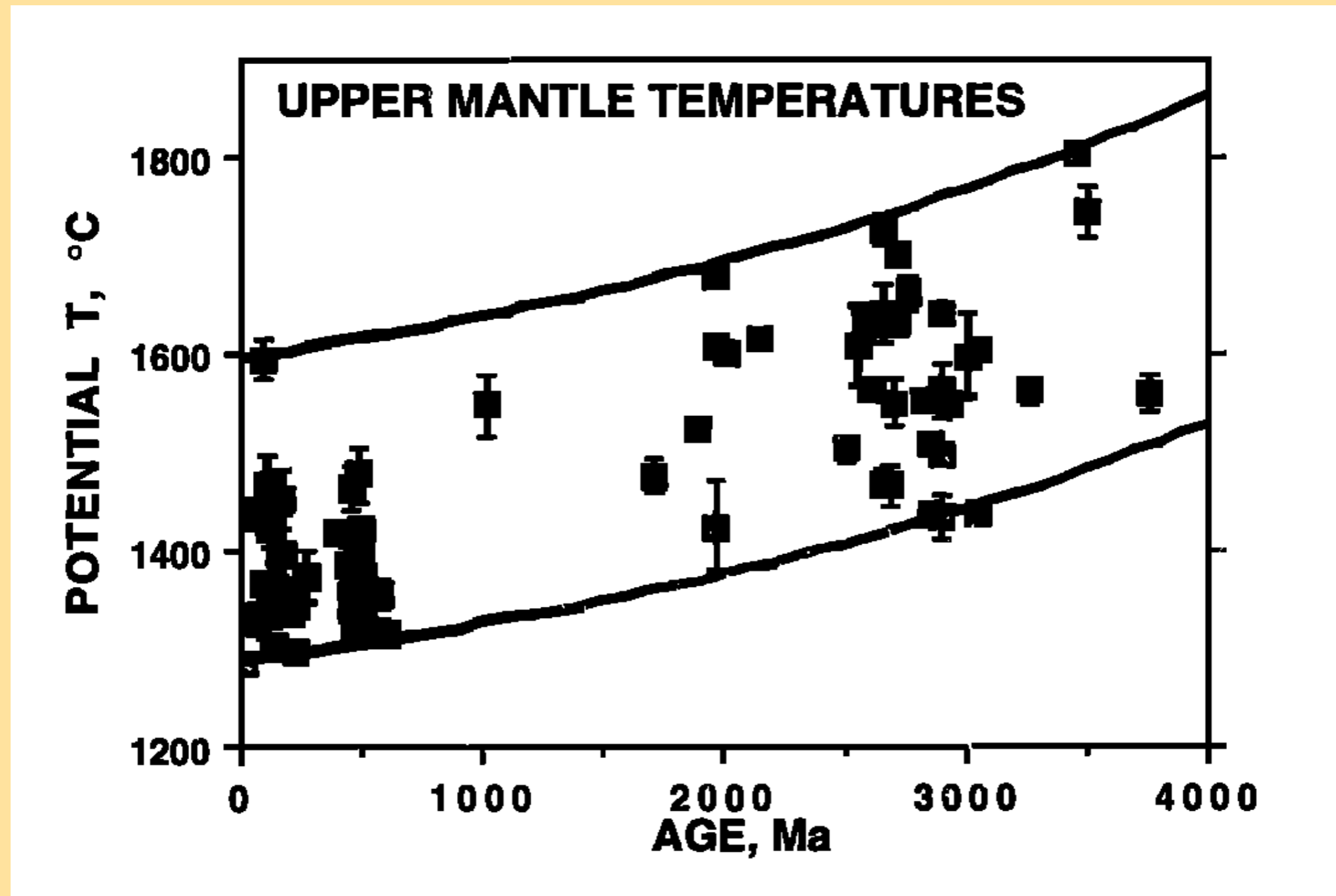


Predicted ^4He :
600 Mmol/yr
Observed:
100 Mmol/yr

Argon

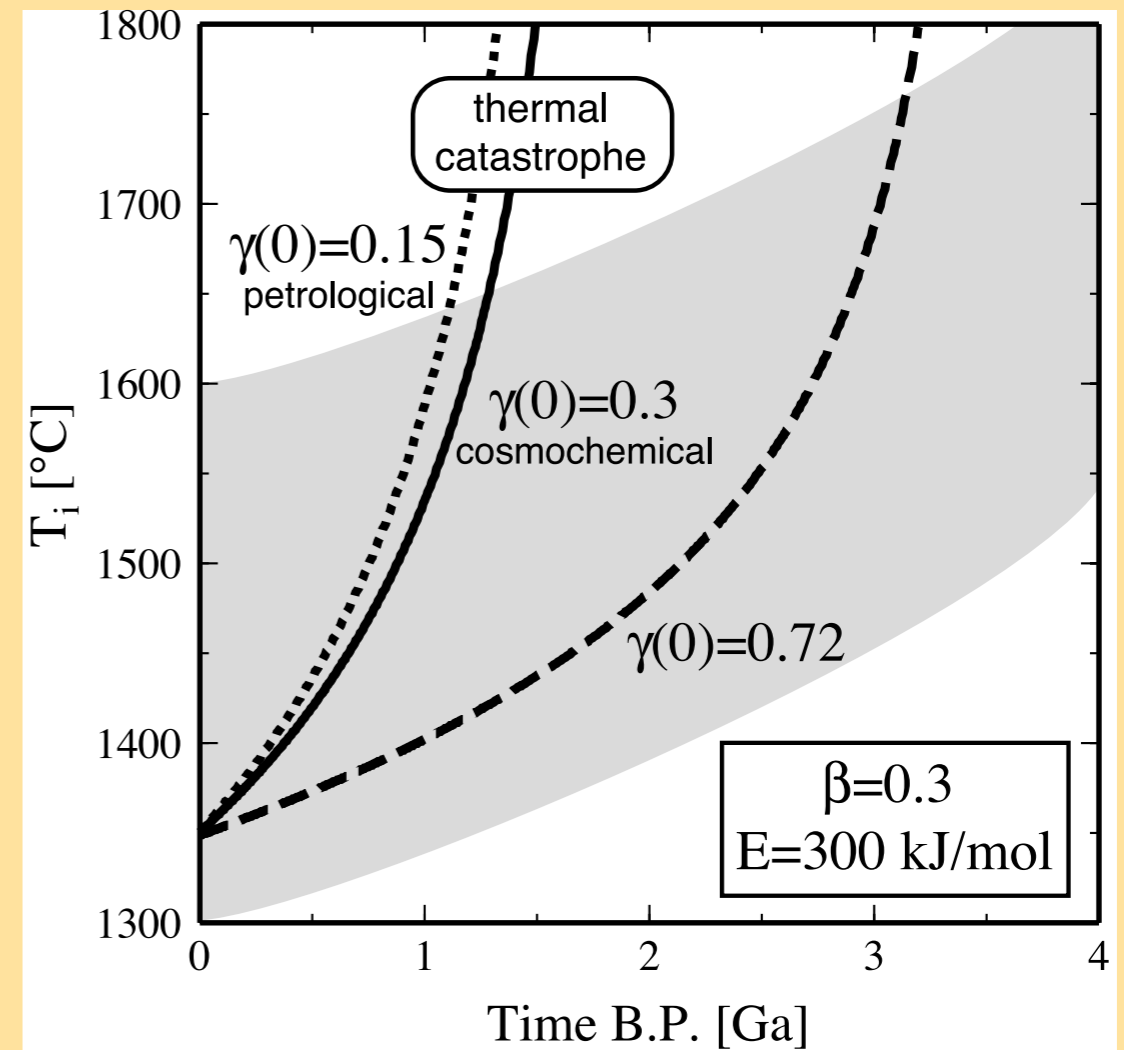
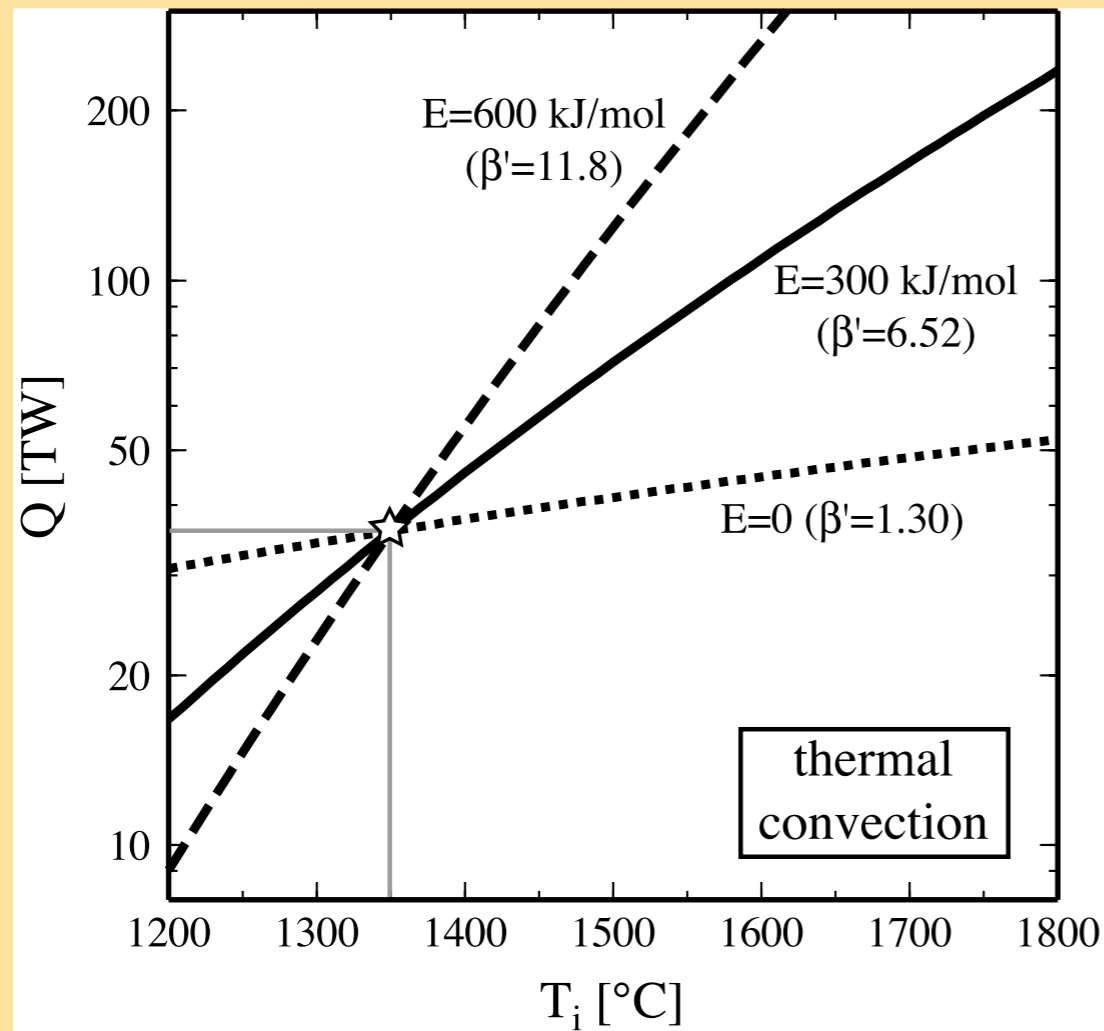
- $^{40}\text{K} \rightarrow ^{40}\text{Ar}$
- ^{40}Ar does not escape from the atmosphere
- Only about half the predicted ^{40}Ar is found in the atmosphere
- The rest must be residing *somewhere*

cooling history



Abbott et al., JGR 1994

The heat flux problem (2)



Ways to limit the heat flux

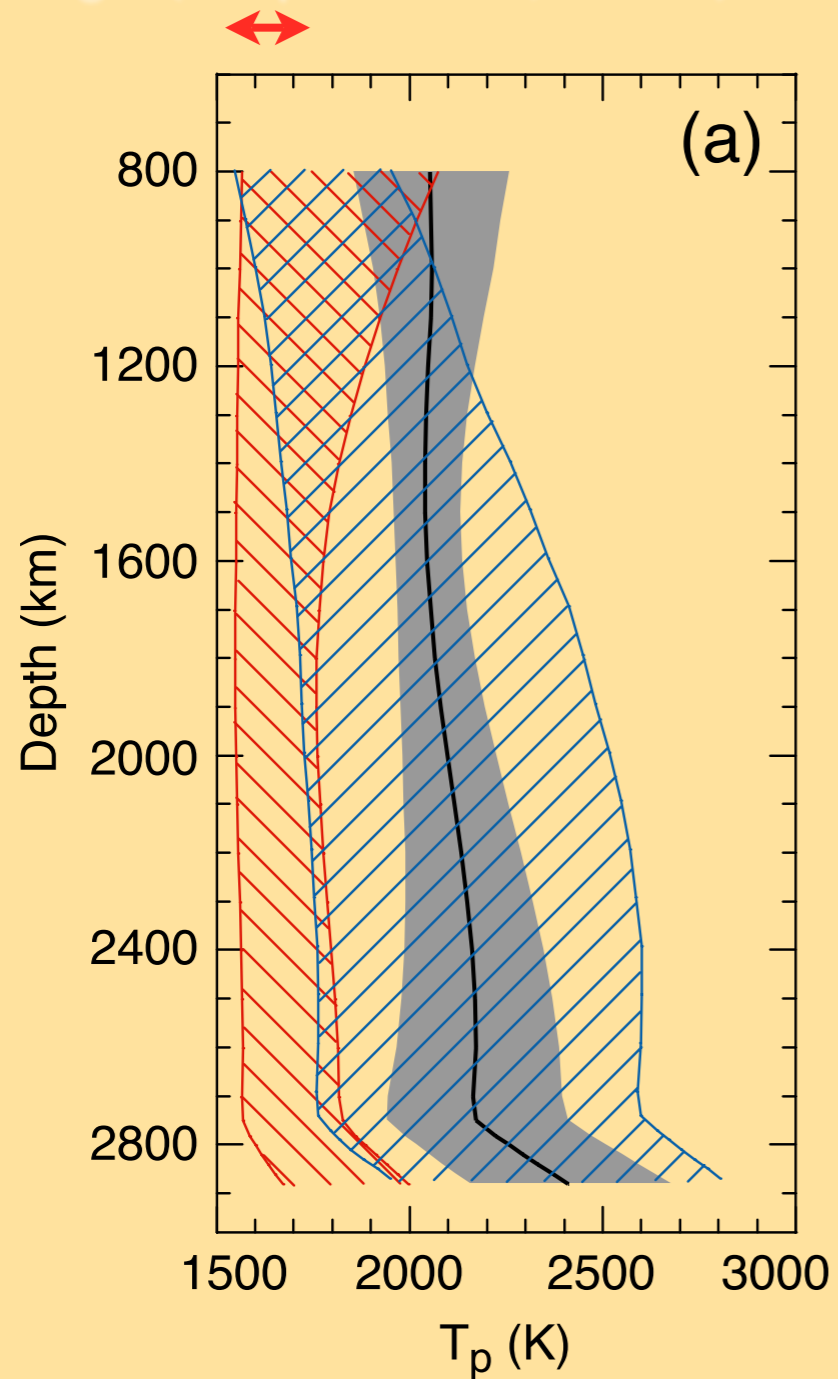
- Plates resist convection (Conrad & Hager, 1999, Korenaga 2006)
- Q is not a simple function of T but depends also on plate configuration (Labrosse & Jaupart, 2007)
- Upper and lower mantle convect separately (McKenzie & Richter, 1981)

Ways to limit the heat flux

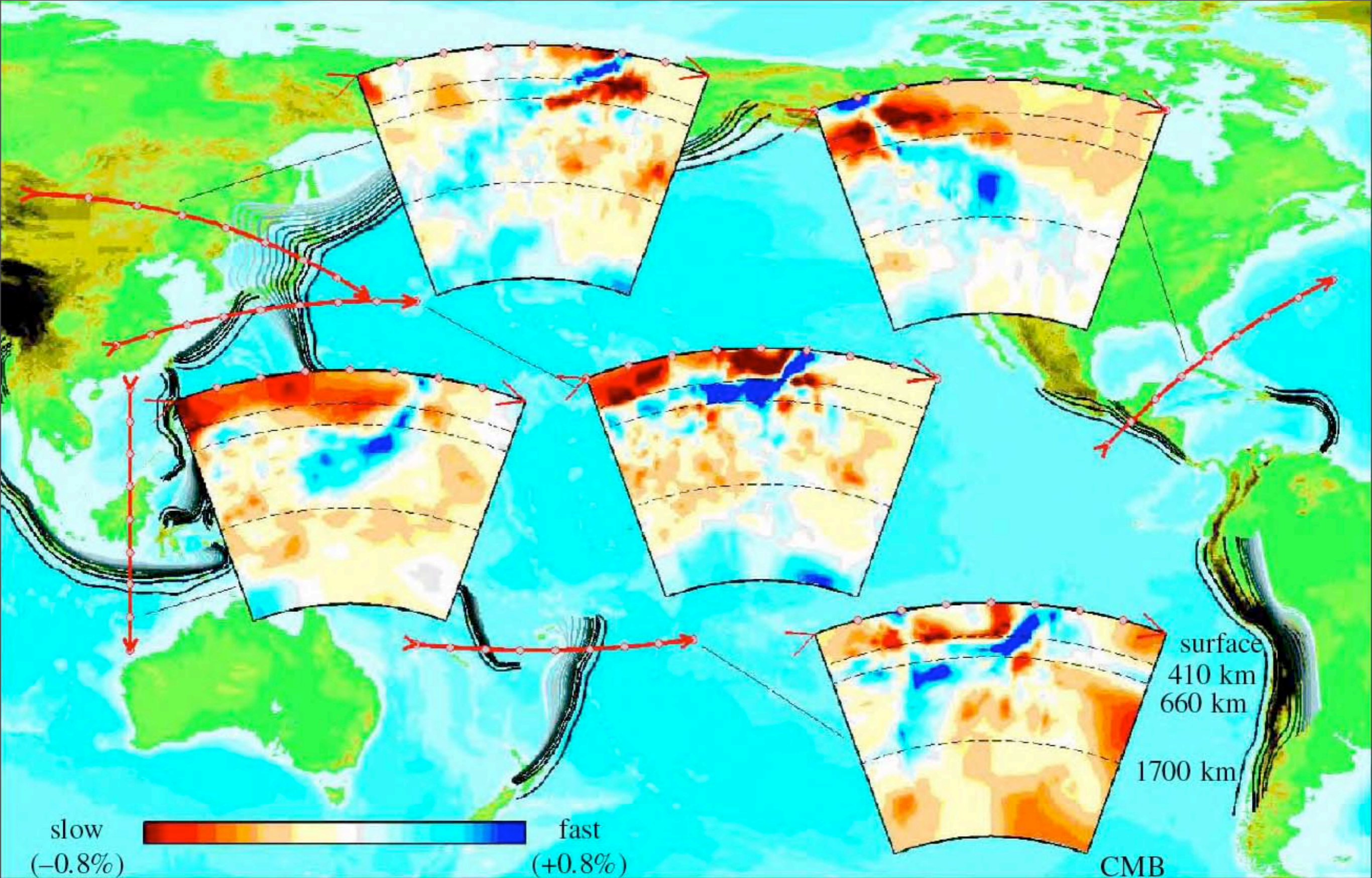
- Plates resist convection (Conrad & Hager, 1999, Korenaga 2006)
- Q is not a simple function of T but depends also on plate configuration (Labrosse & Jaupart, 2007)
- Upper and lower mantle convect separately (McKenzie & Richter, 1981) **This would imply a thermal boundary layer (TBL) at 660 km.**

TBL: what does PREM tell us?

Upper mantle potential temperature range (Jaupart et al, 2008)

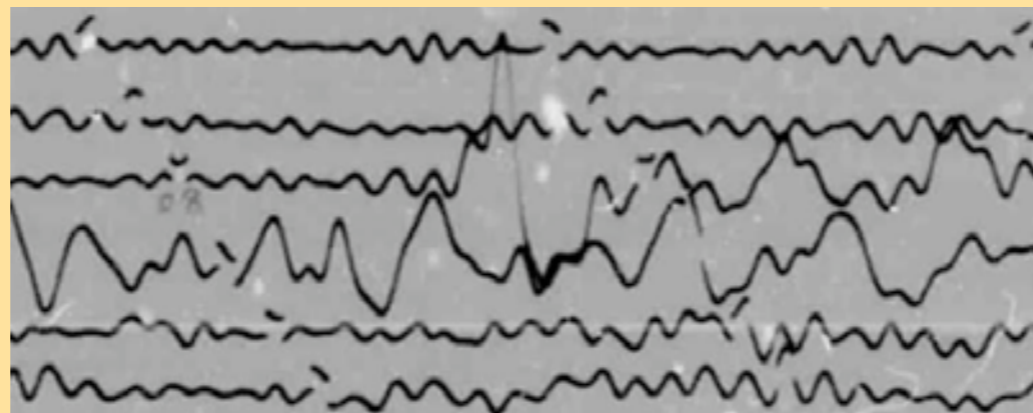
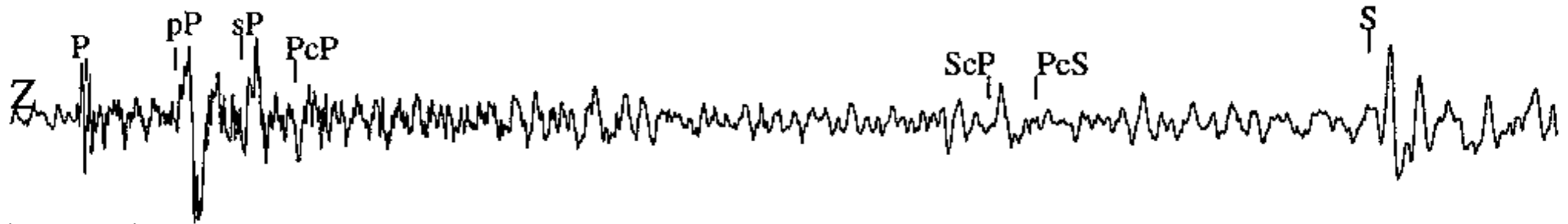


Deschamps & Trampert, EPSL 2004



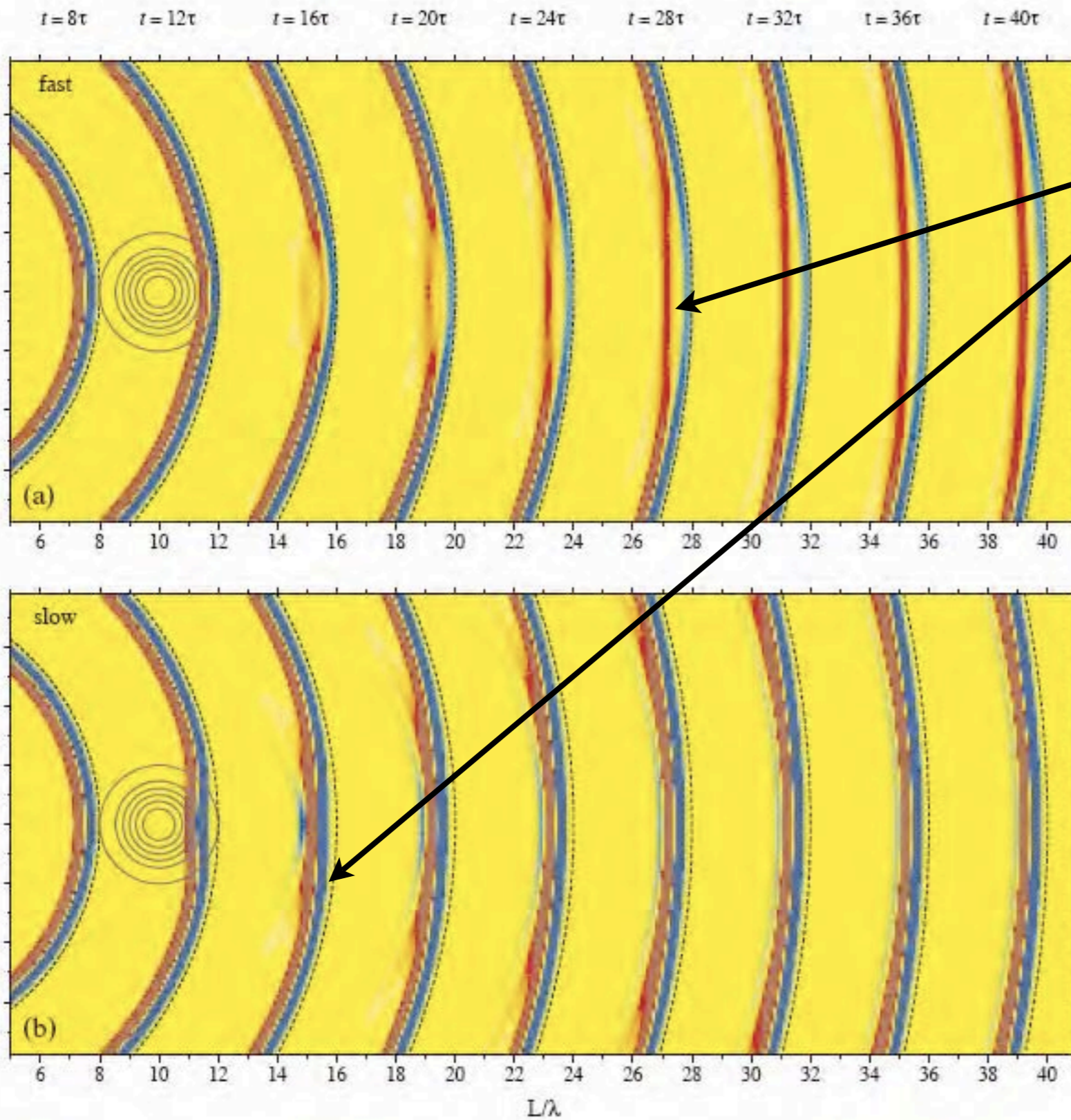
Albarède & Van der Hilst, 2002

Intermezzo: from onset times to cross-correlations

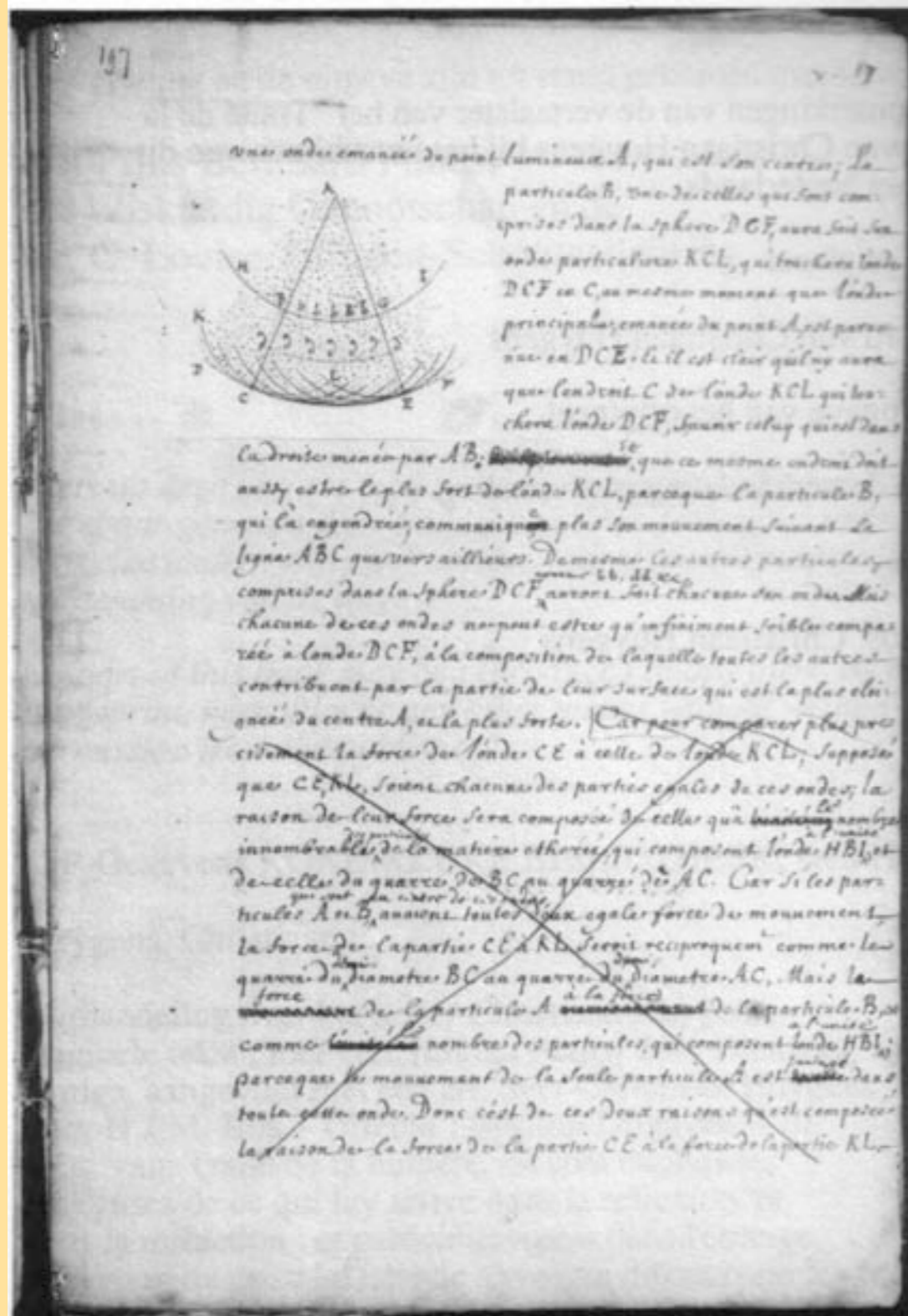


60 s

The most dramatic effects are at later times in the waveform!

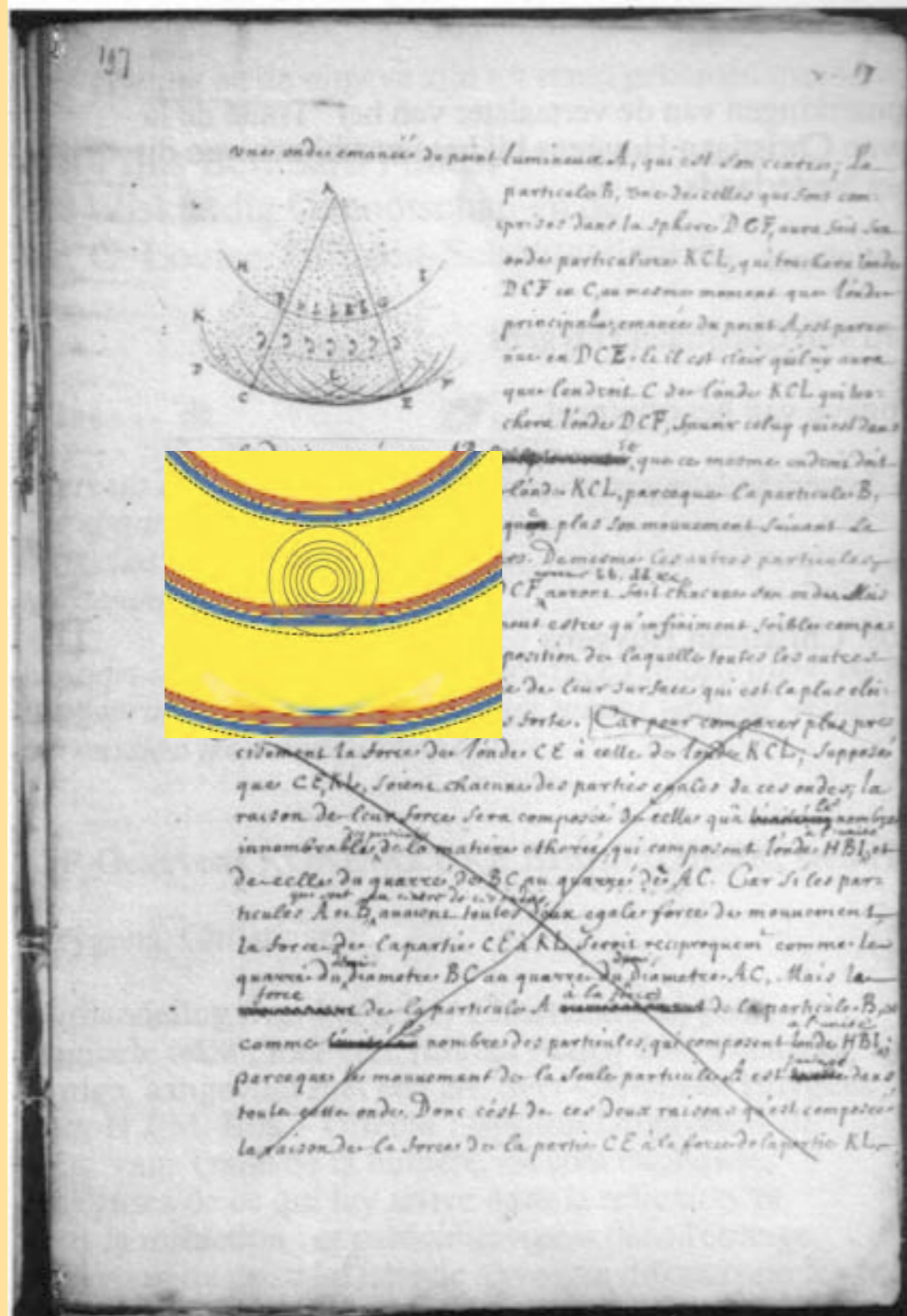


Wave diffraction



Christiaan Huygens (1629-1696)

Wave diffraction

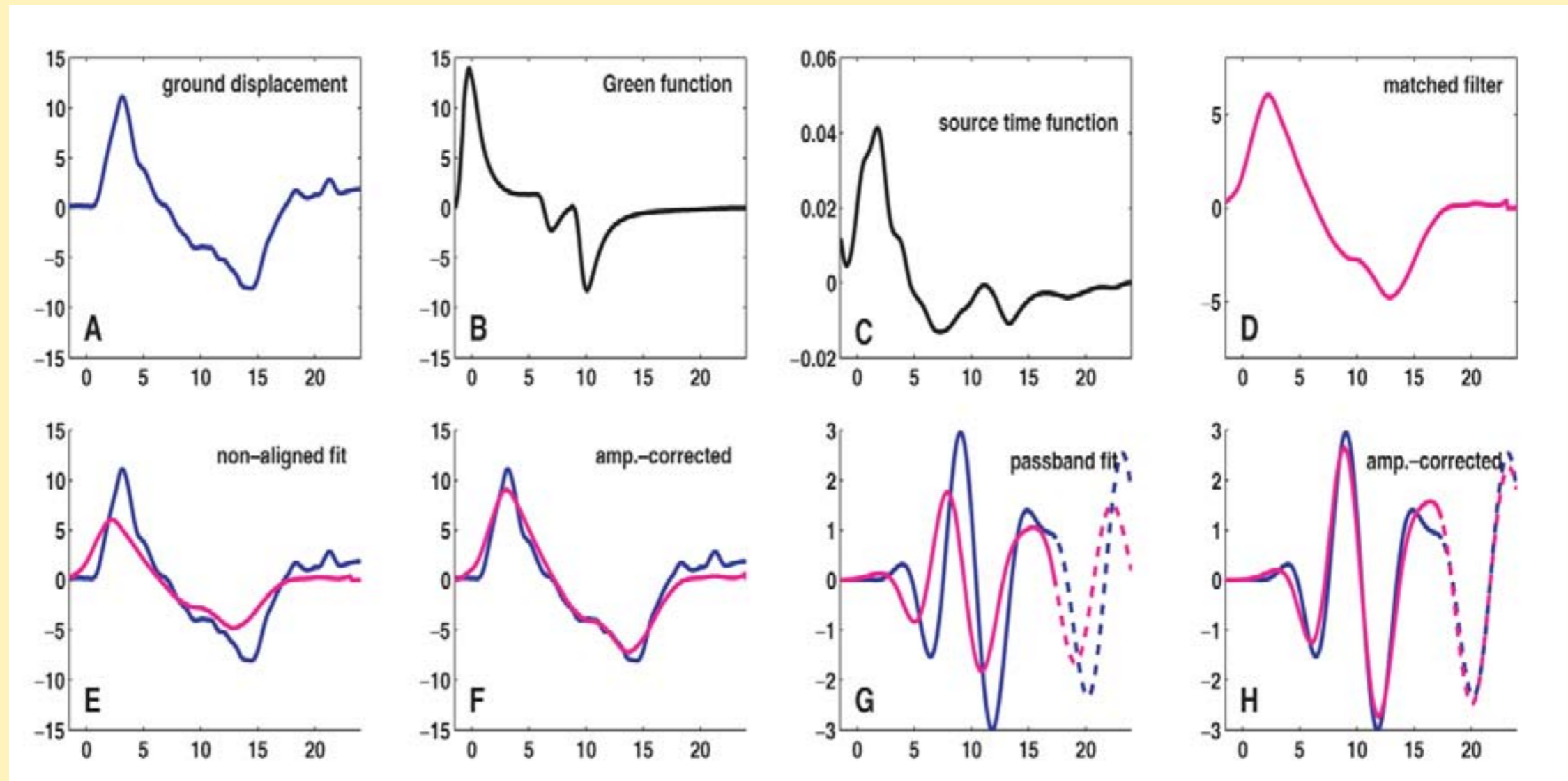


Christiaan Huygens (1629-1696)

Waves may take detours



Body wave cross-correlations

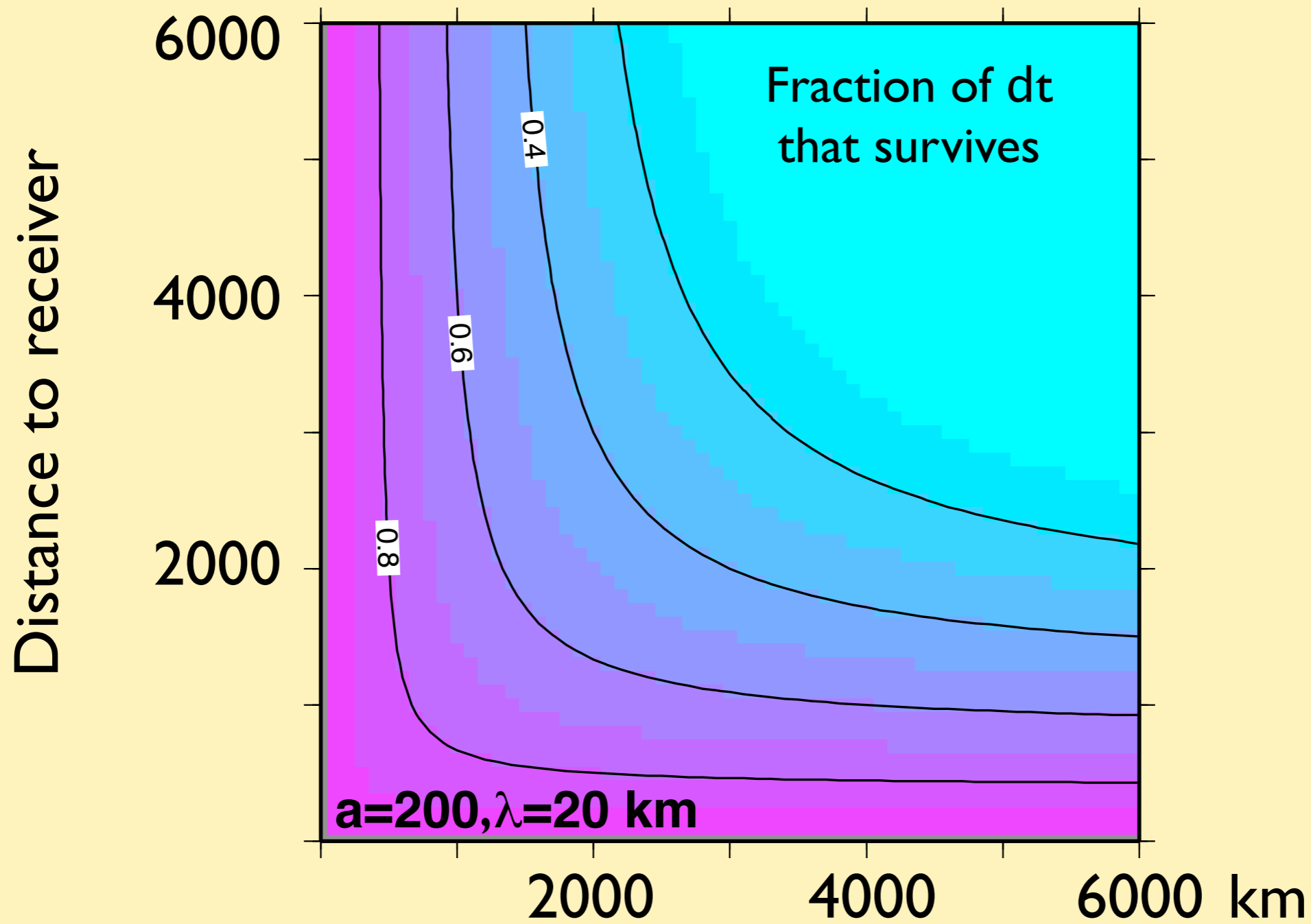


Cross-correlation

$$\gamma(t) = \int s(t')u(t' - t)dt'$$

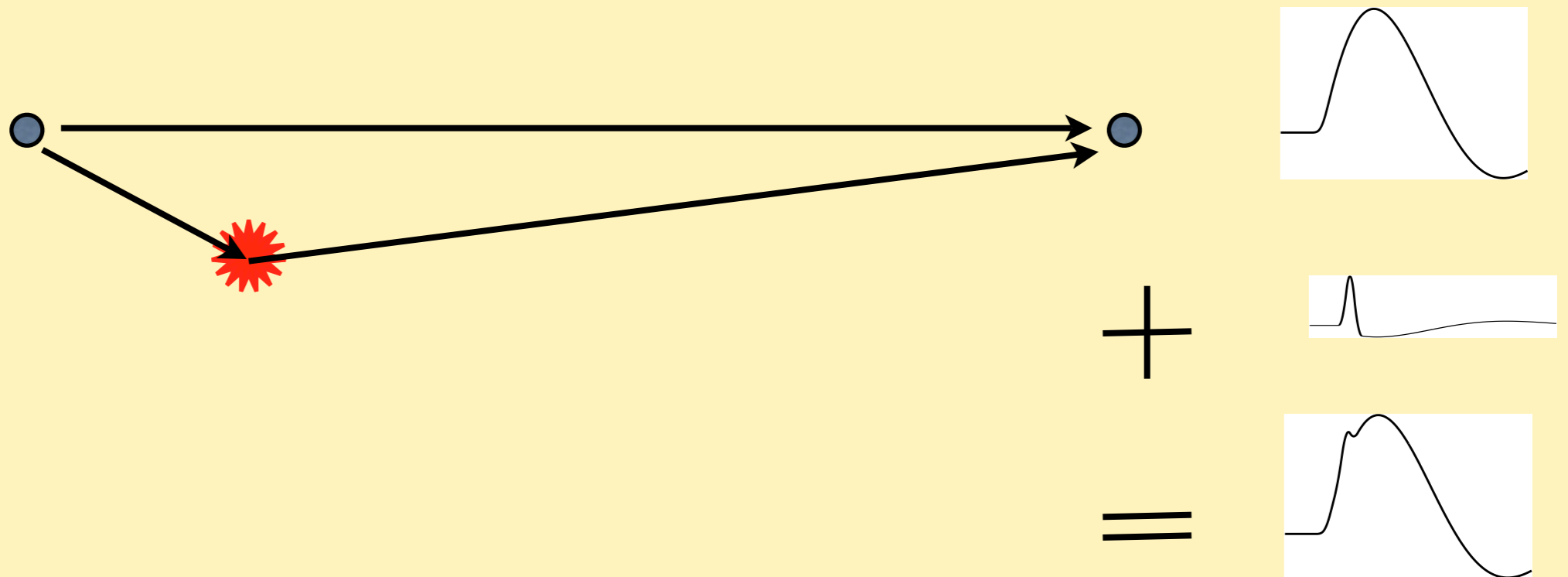
$$\sigma_{\text{CRLB}}^2 = \frac{3}{8\pi^2} \frac{1 + 2\text{SNR}}{\text{SNR}^2} \frac{1}{\Delta f^3 T_w}$$

Healing of cross-correlation delays (Period=2 s)



*Based on numerical simulations
by Hung et al., GJI 2001.*

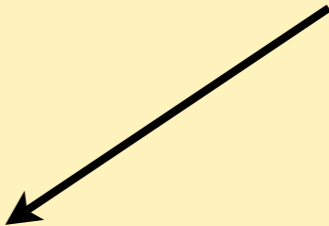
Born theory = first order
perturbation *of an early arrival*



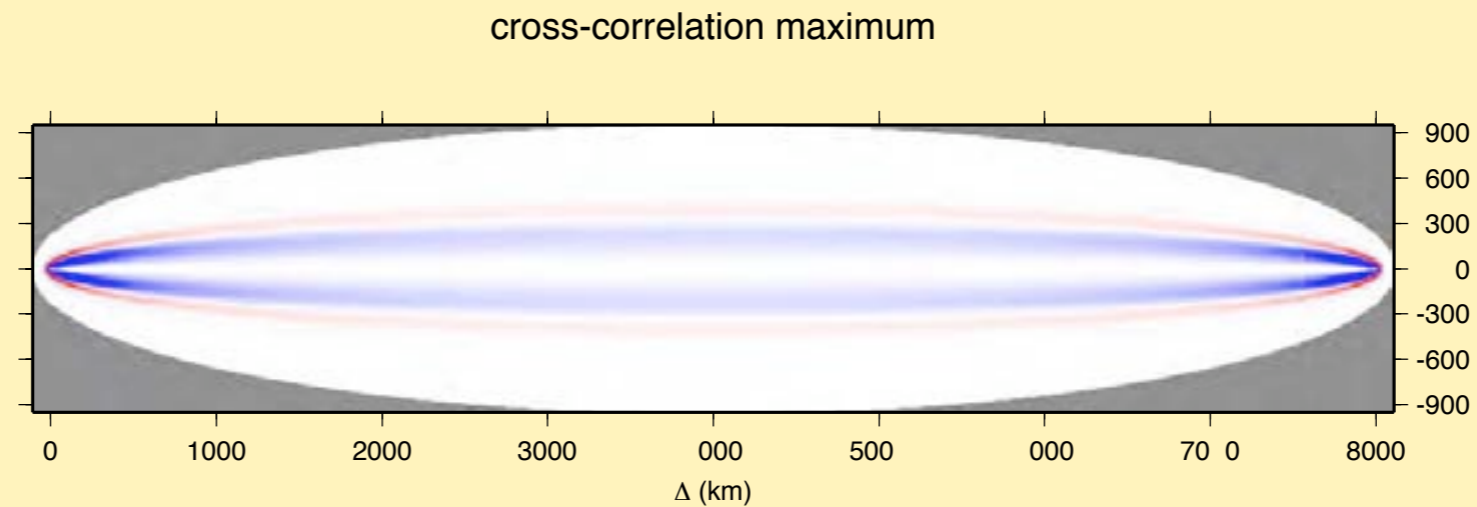
Scattered wave perturbs crosscorrelation

$$\delta T = -\frac{\delta \dot{\gamma}(0)}{\ddot{\gamma}(0)} = -\frac{\int_{-\infty}^{\infty} \dot{u}(t') \delta u(t') dt'}{\int_{-\infty}^{\infty} \ddot{u}(t') u(t') dt'} .$$

Born



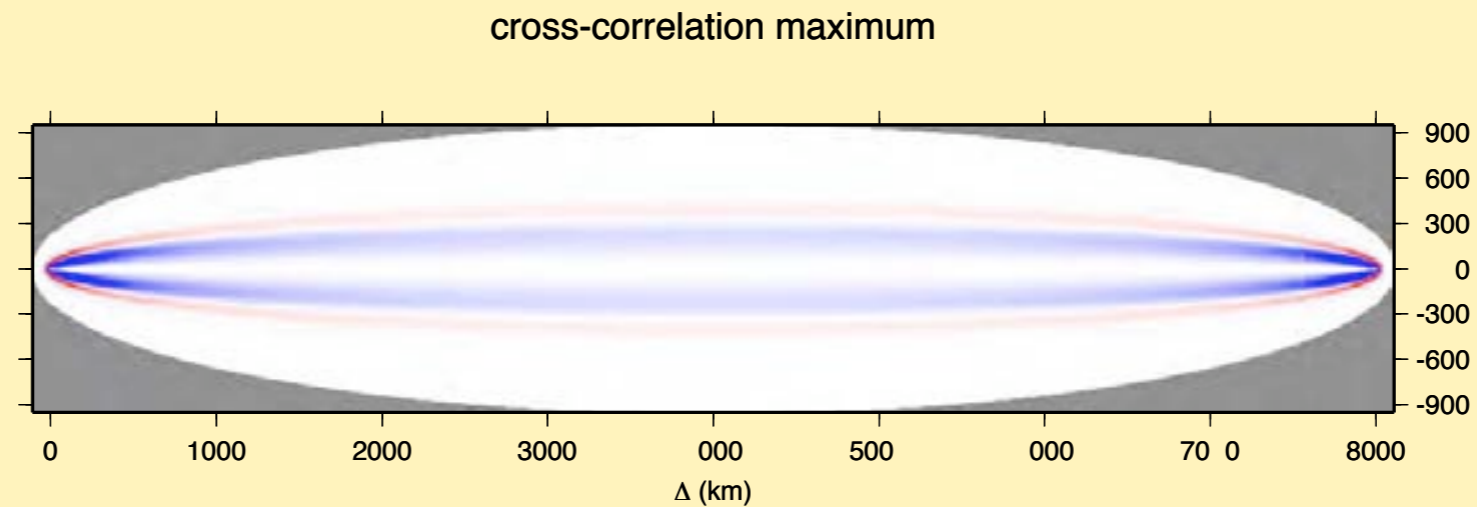
Scattered waves + delay from cross-correlation \rightarrow “banana-doughnut” kernels



$$\delta T = \int K_P \left(\frac{\delta V_P}{V_P} \right) d^3 r_x$$

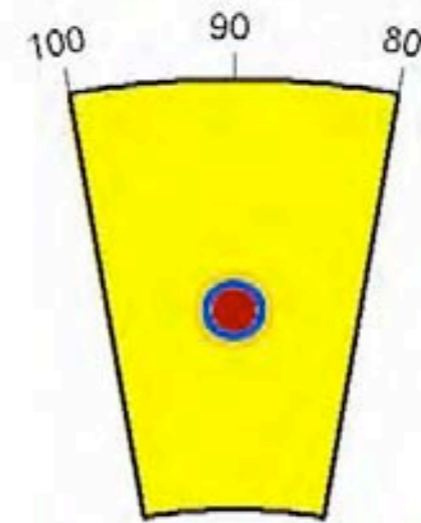
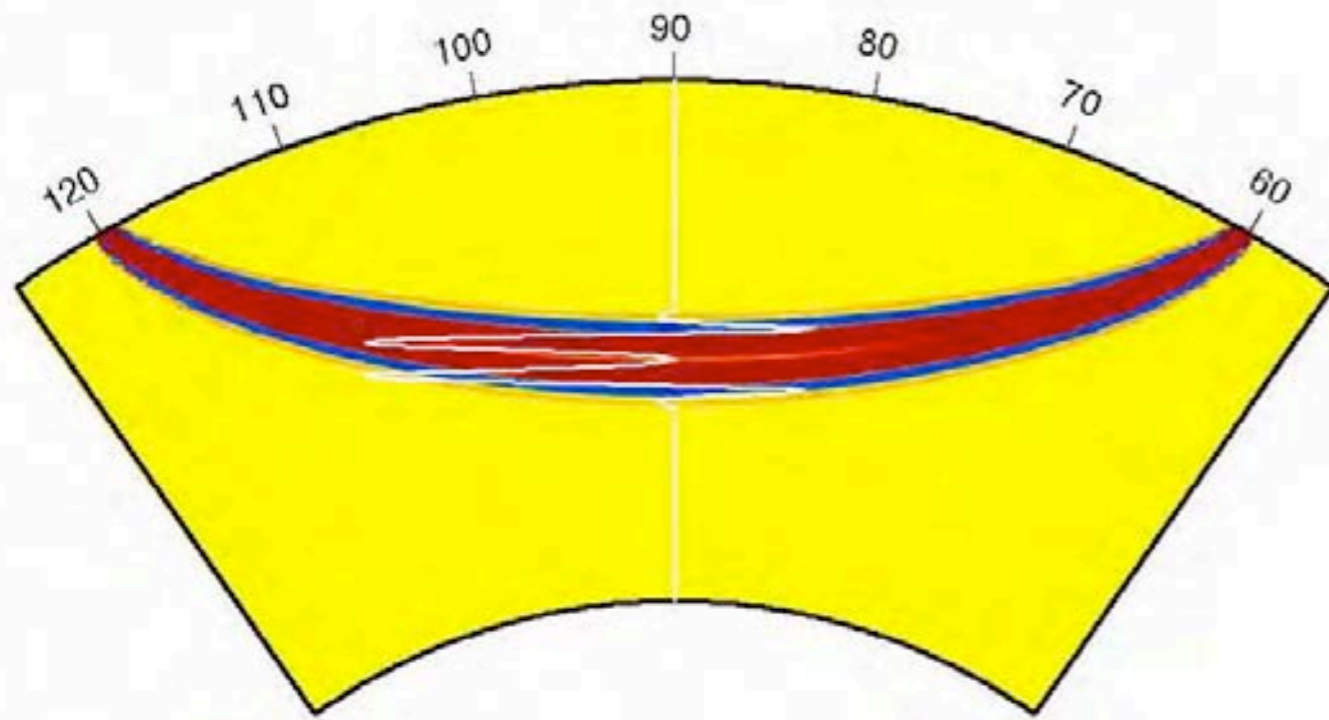
$$= \frac{\int_{-\infty}^{\infty} \dot{u}(t') \delta u(t') dt'}{\int_{-\infty}^{\infty} \ddot{u}(t') u(t') dt'}$$

Scattered waves + delay from cross-correlation \rightarrow “banana-doughnut” kernels

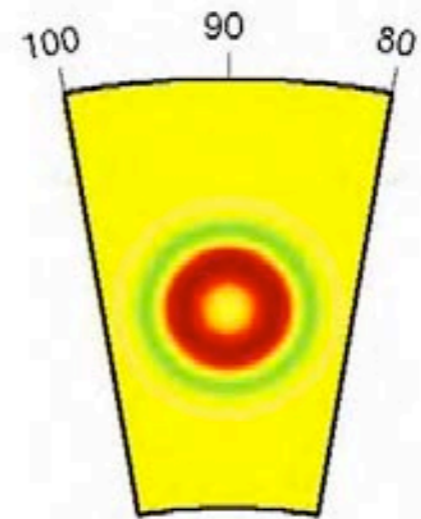
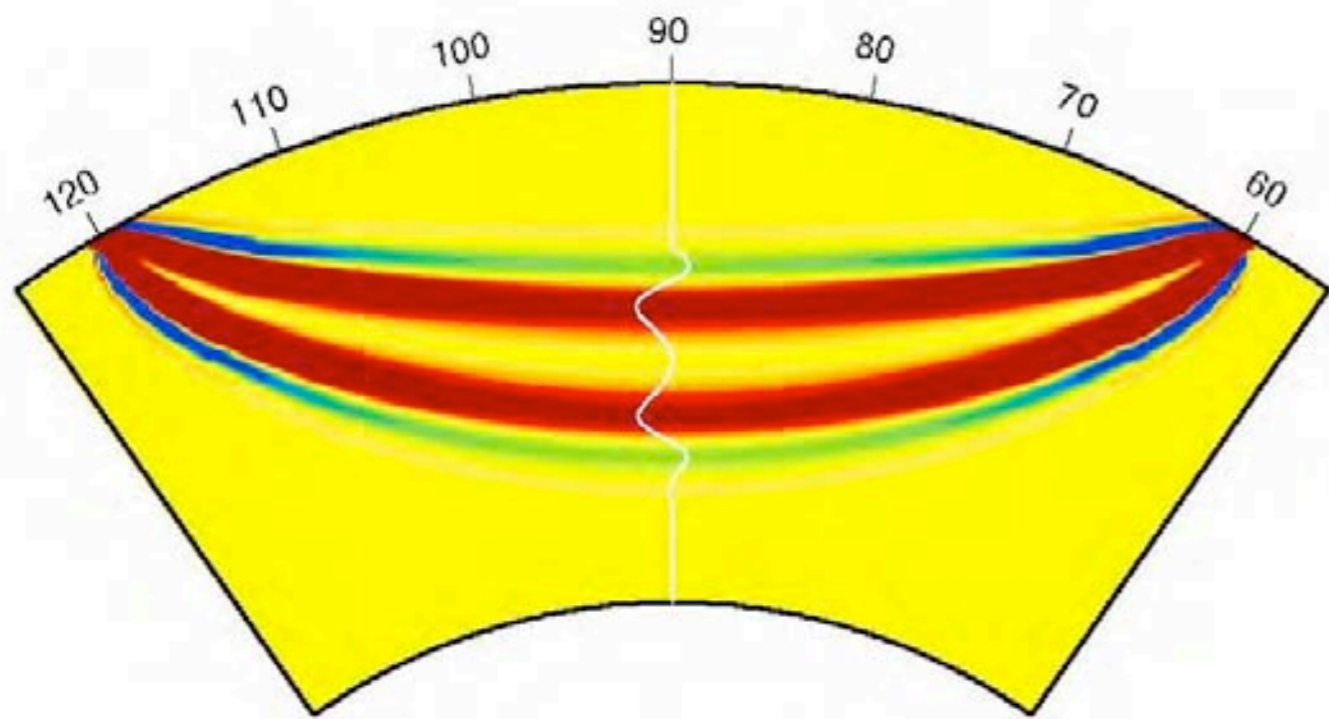


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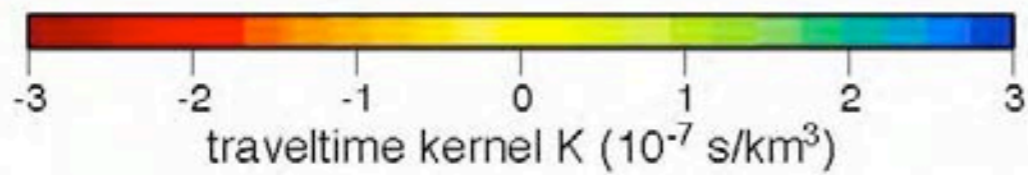
$$= \frac{\int_{-\infty}^{\infty} \dot{u}(t') \delta u(t') dt'}{\int_{-\infty}^{\infty} \ddot{u}(t') u(t') dt'}$$



$T=2s$

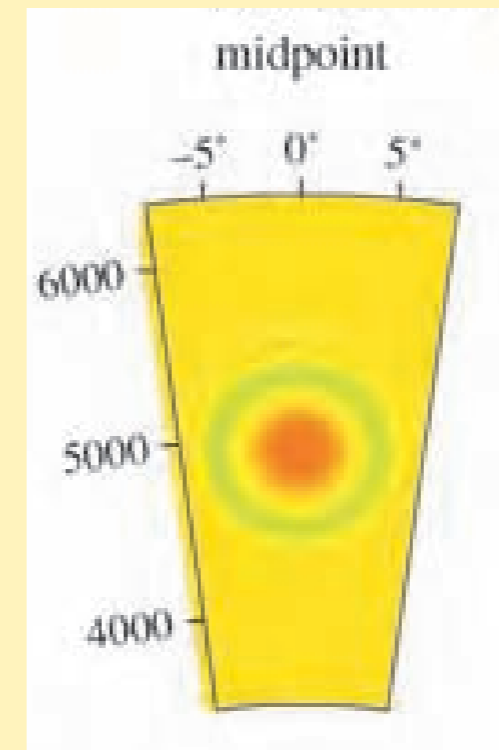
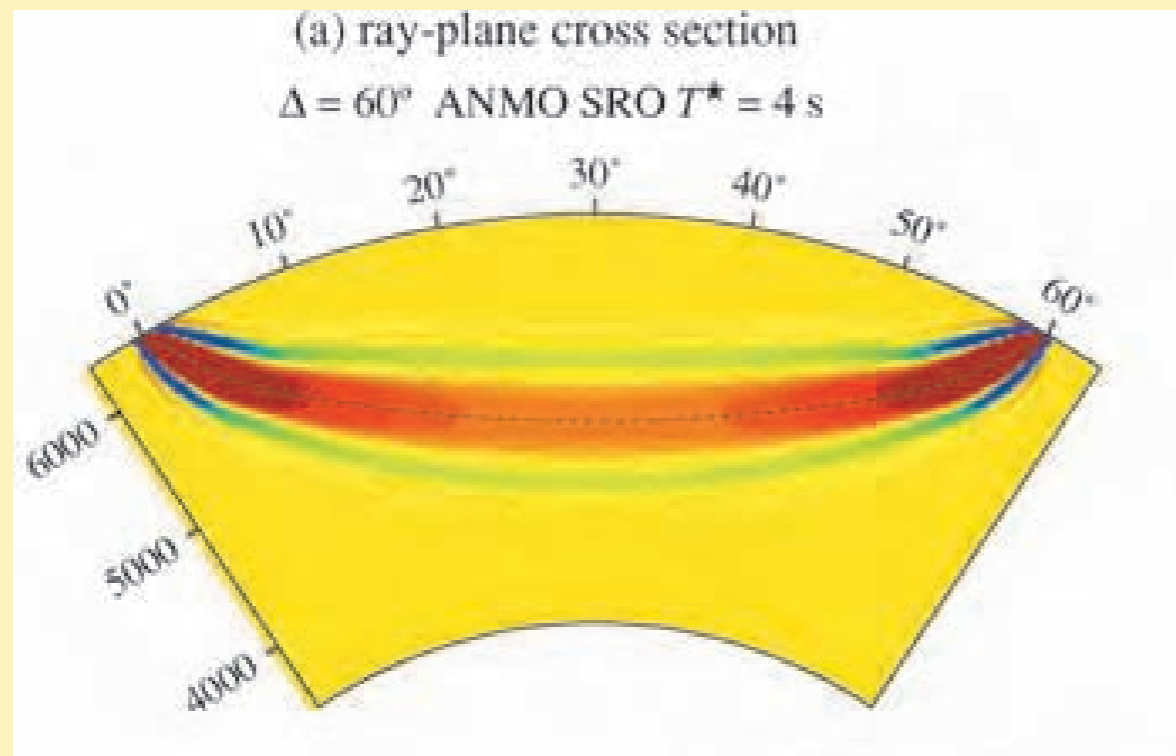


$T=20s$

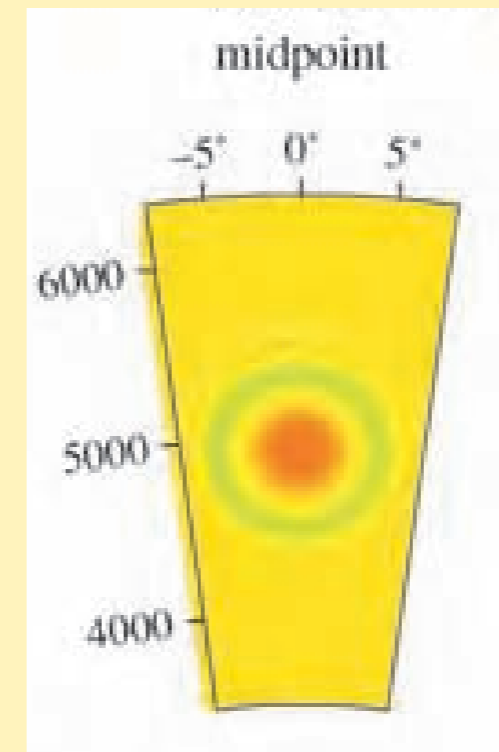
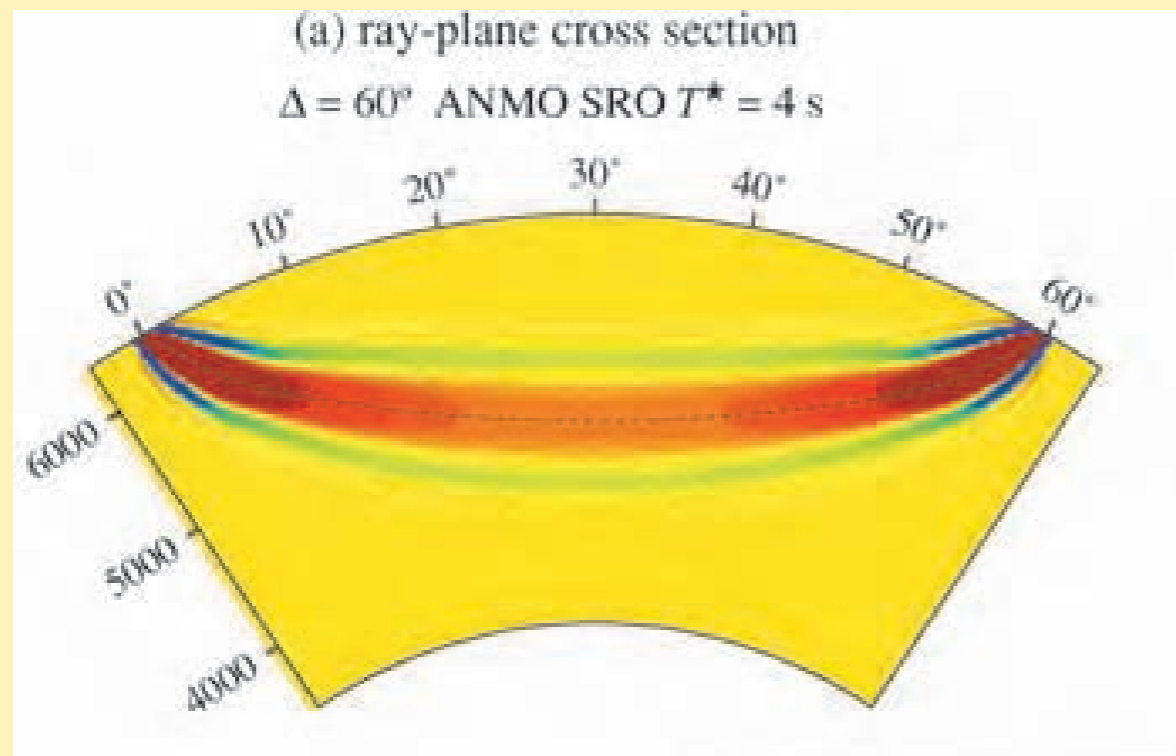


The period T determines the width of the kernel

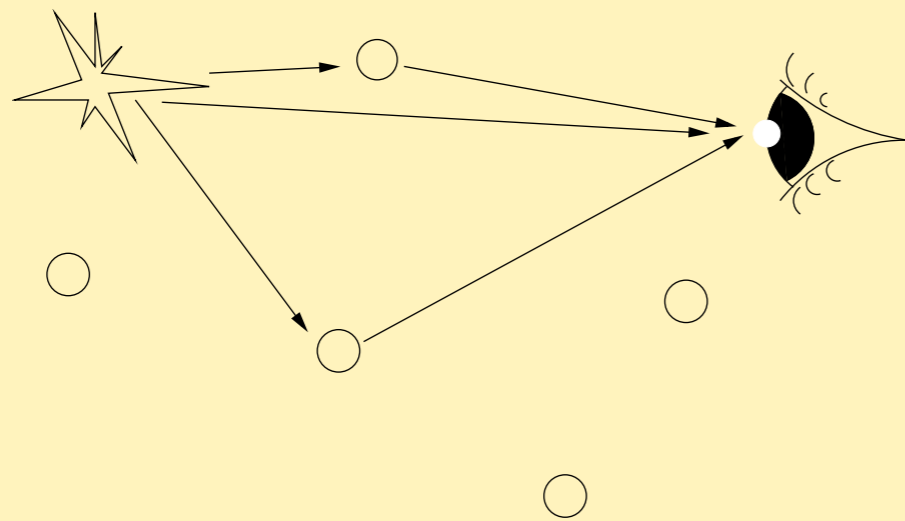
Amplitude kernels



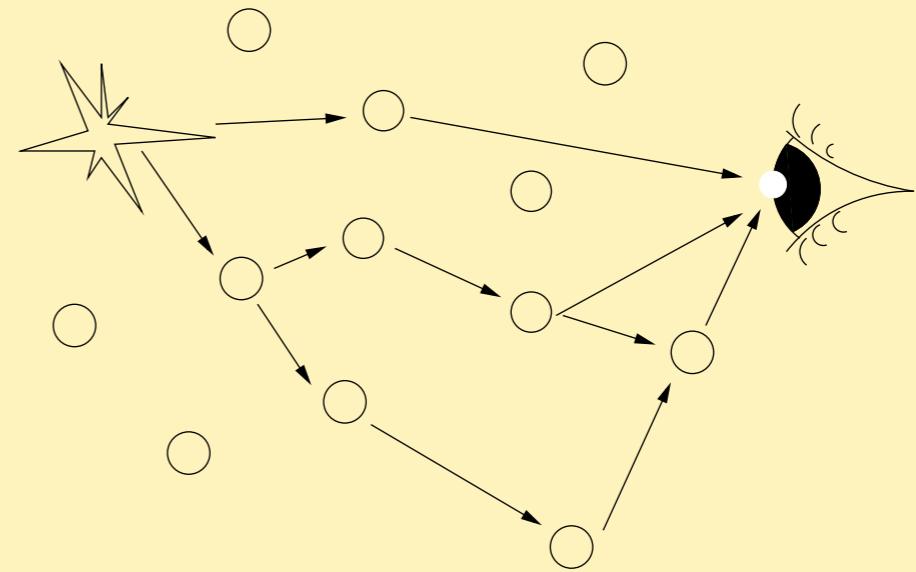
Amplitude kernels



modeling waveforms: Single or multiple scattering?



(a) Diffusion simple



(b) Diffusion multiple

Single=early

Multiple =late

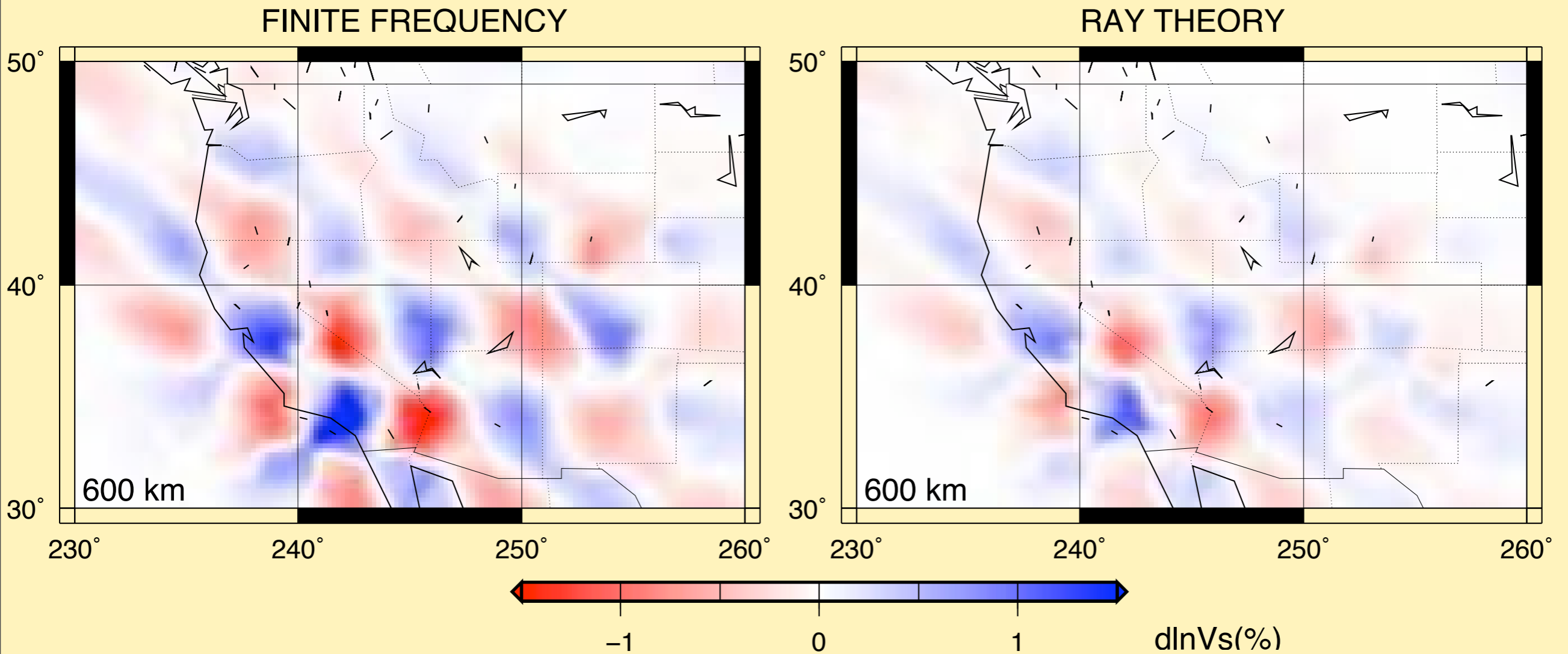
Forward scattering retains information



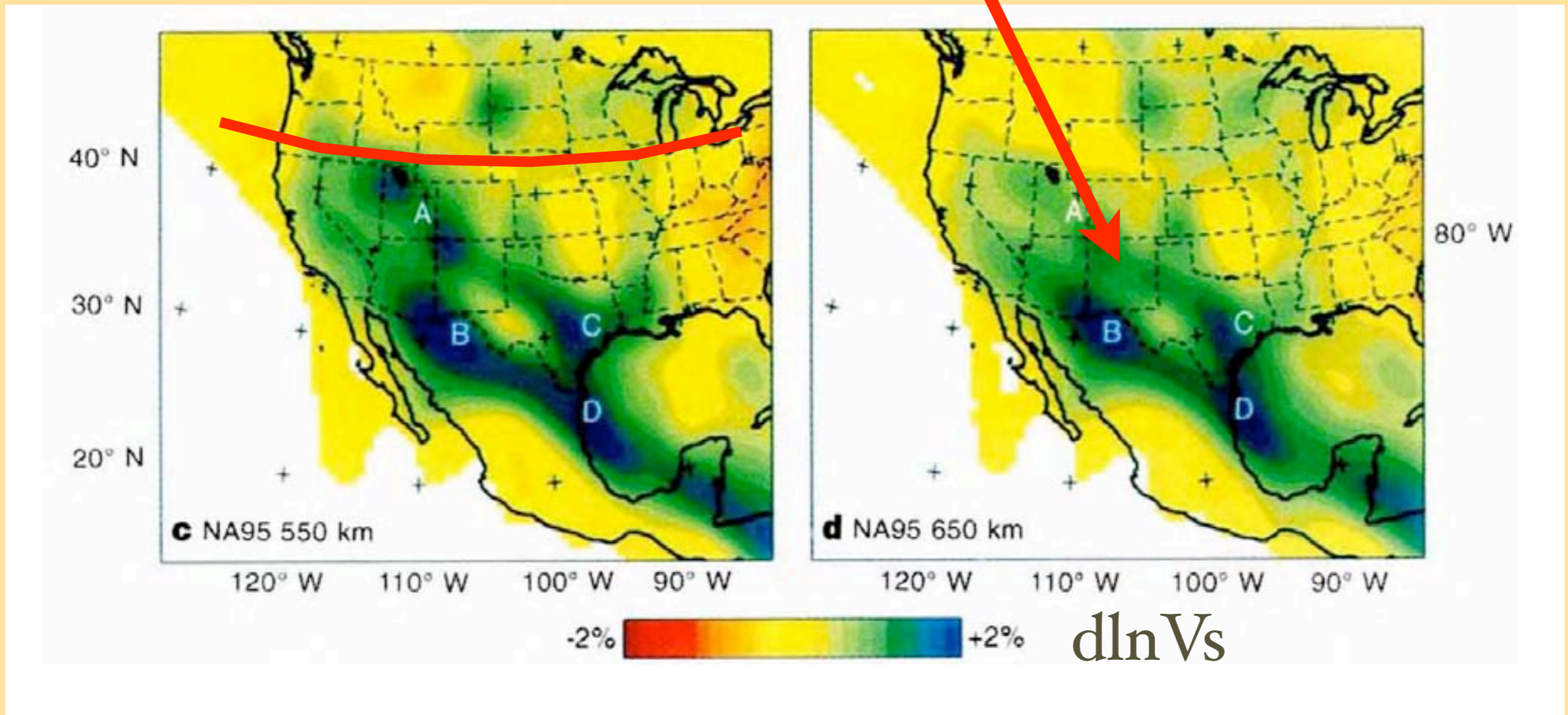
Multiple scattering: ill posed inverse problem



Resolution gain (multiple frequency, delays & amplitudes)

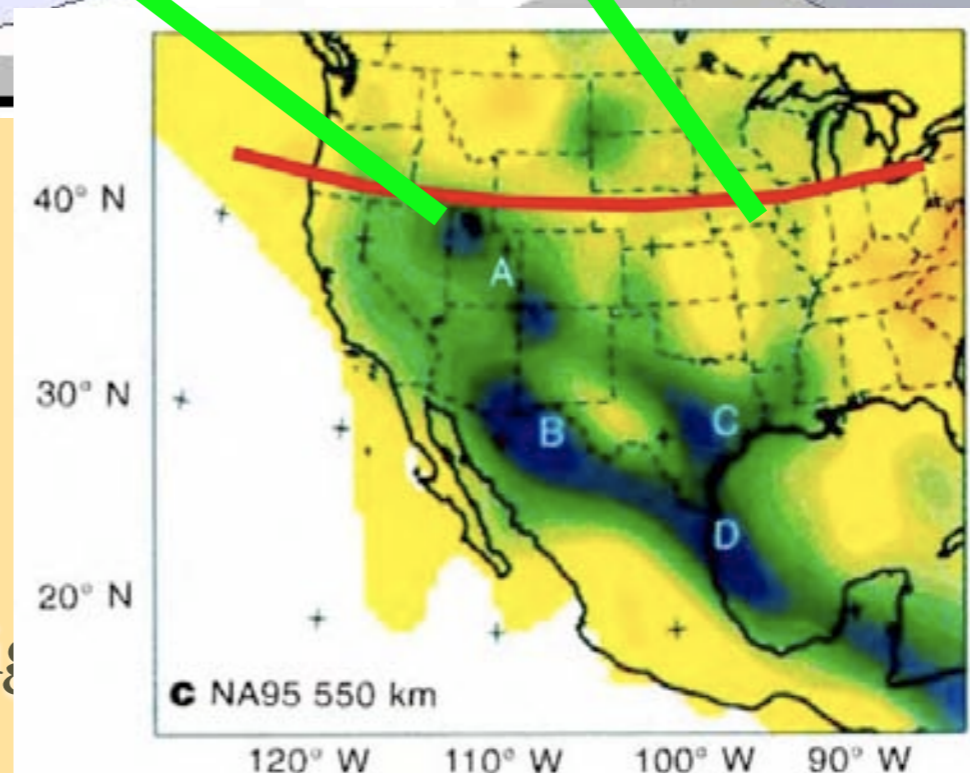
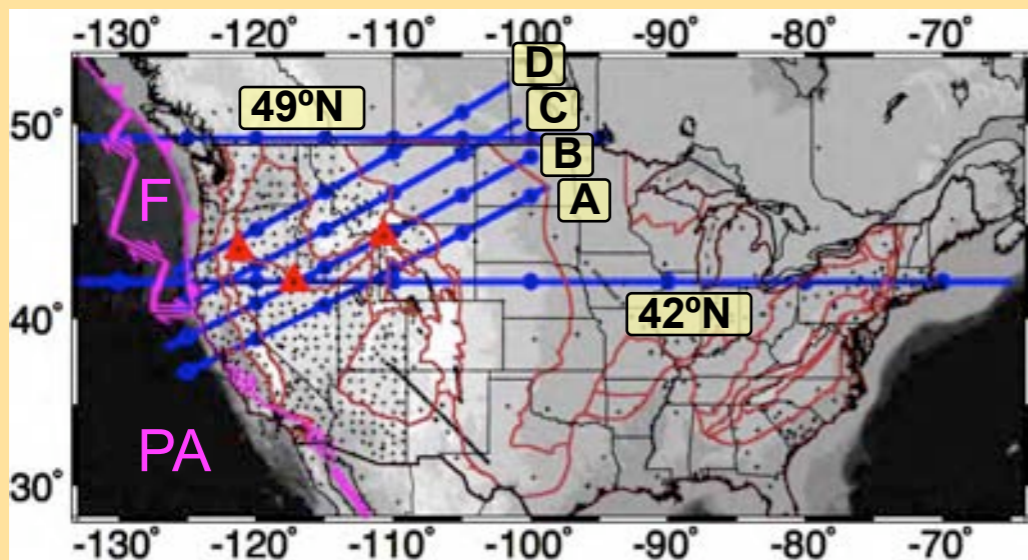
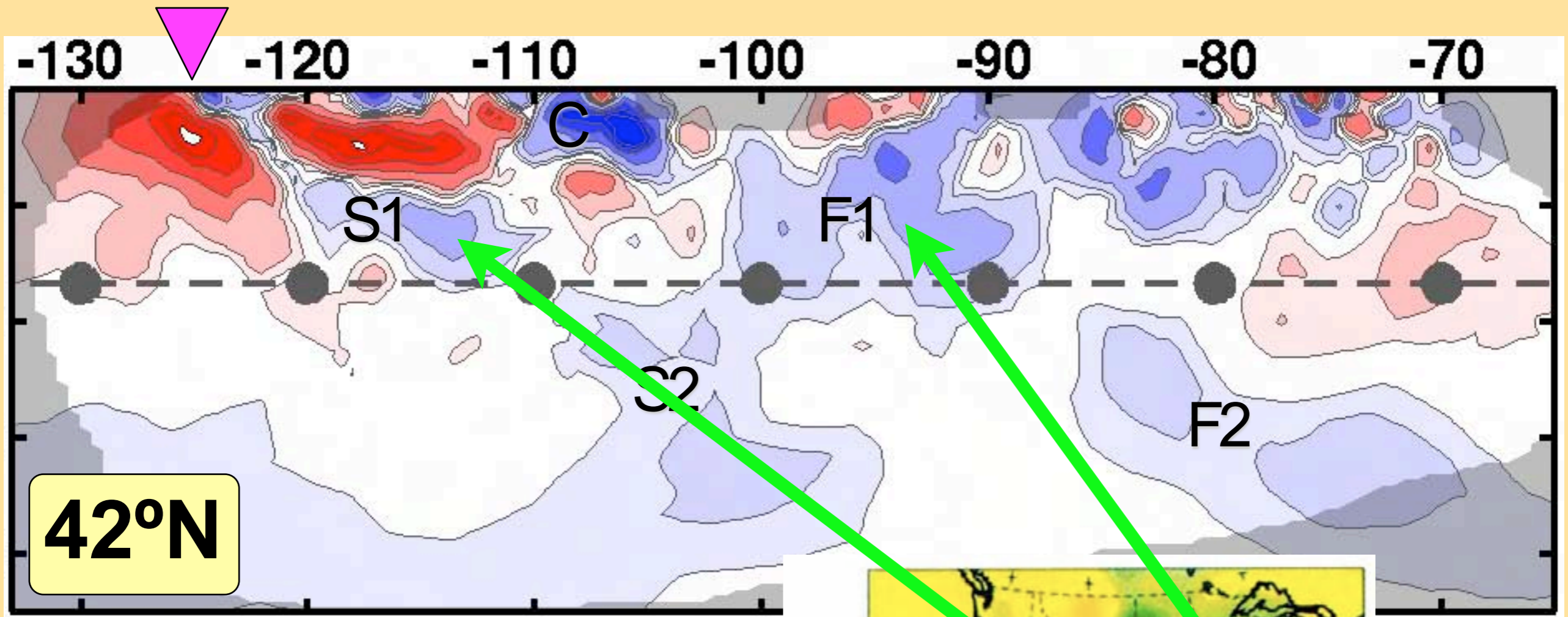


Early surface wave images: Farallon stuck in the TZ

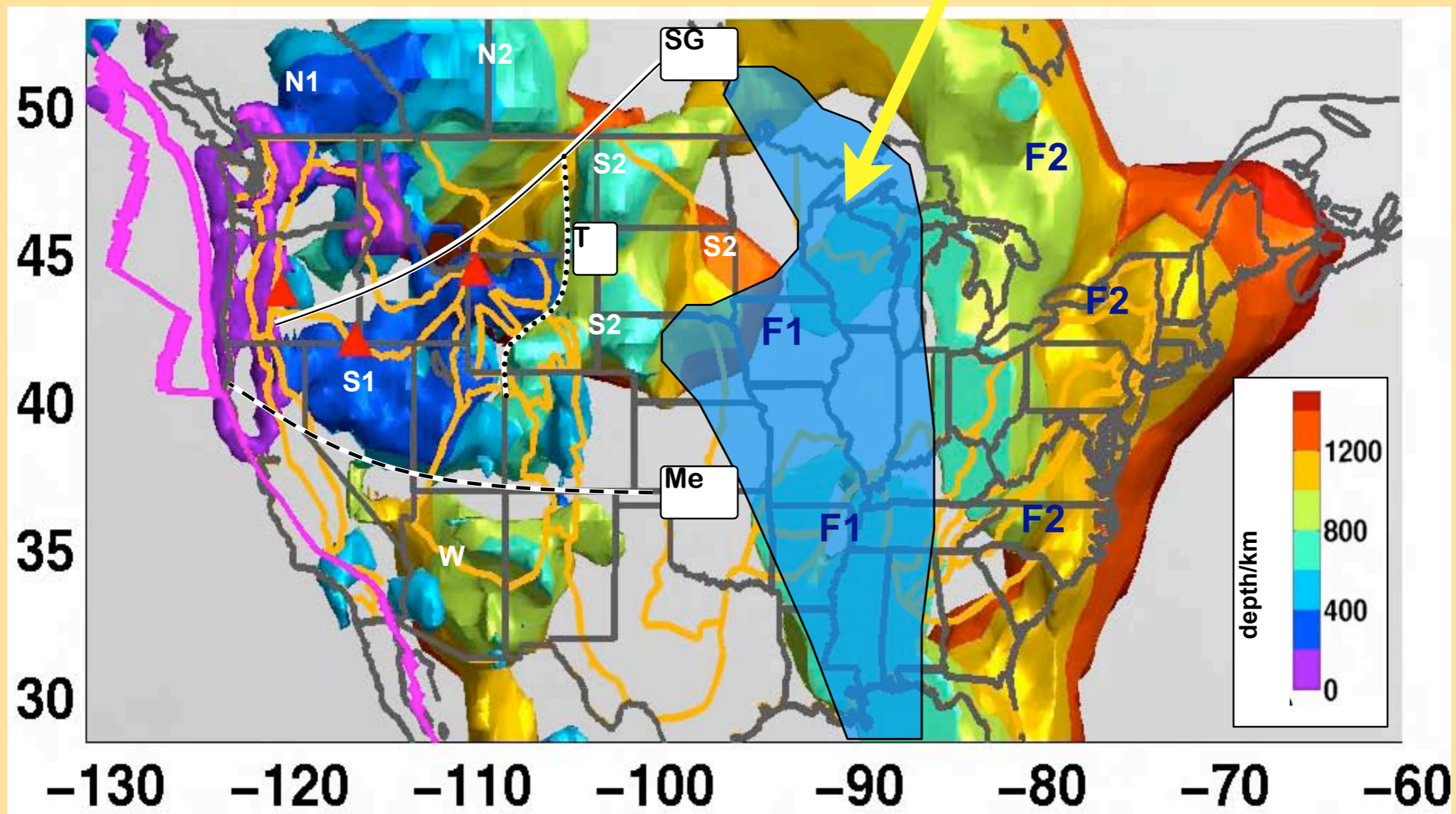


Van der Lee & Nolet, Nature 1997

Finite frequency P wave images and USArray

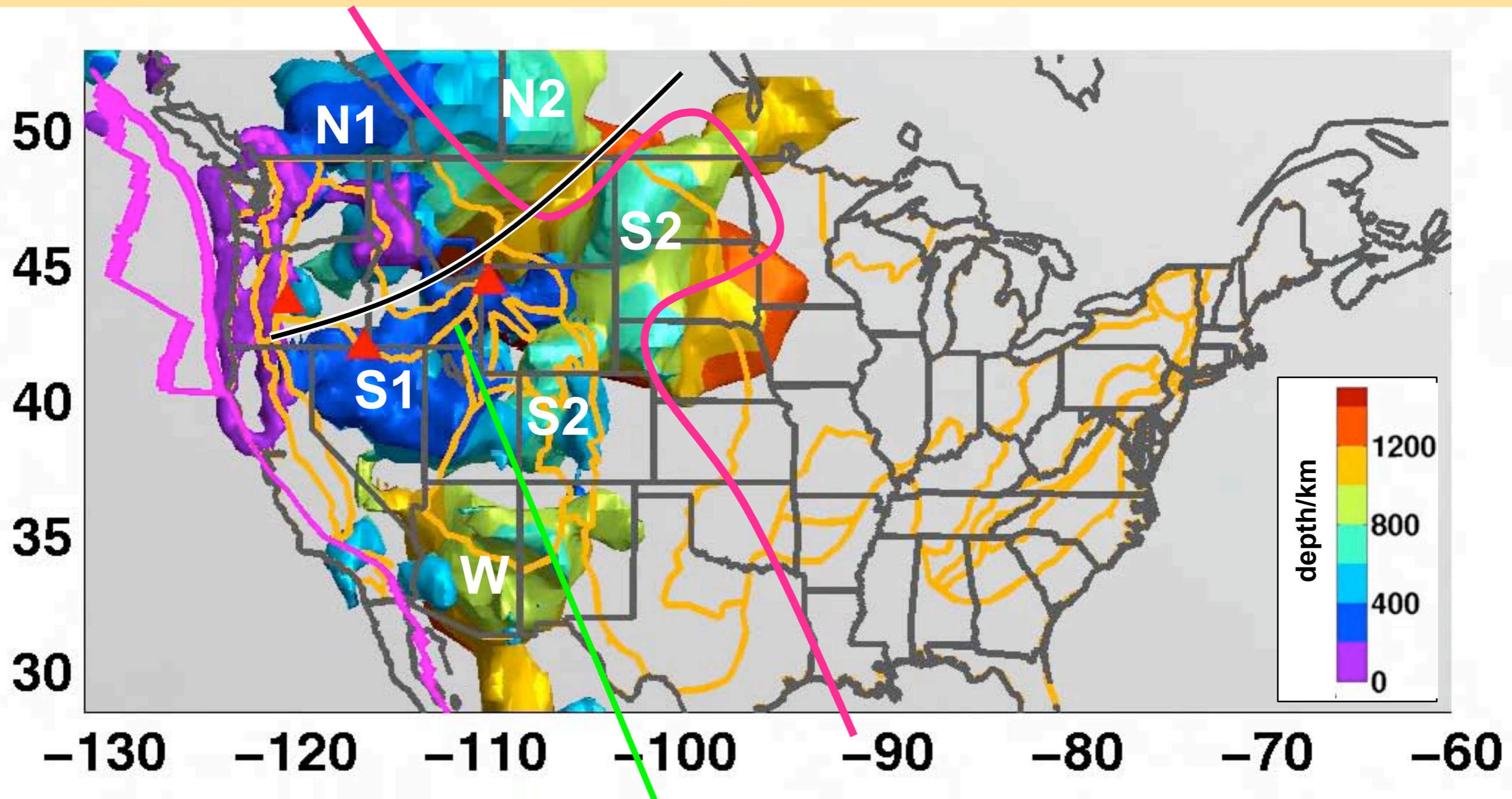


Farallon stuck in TZ

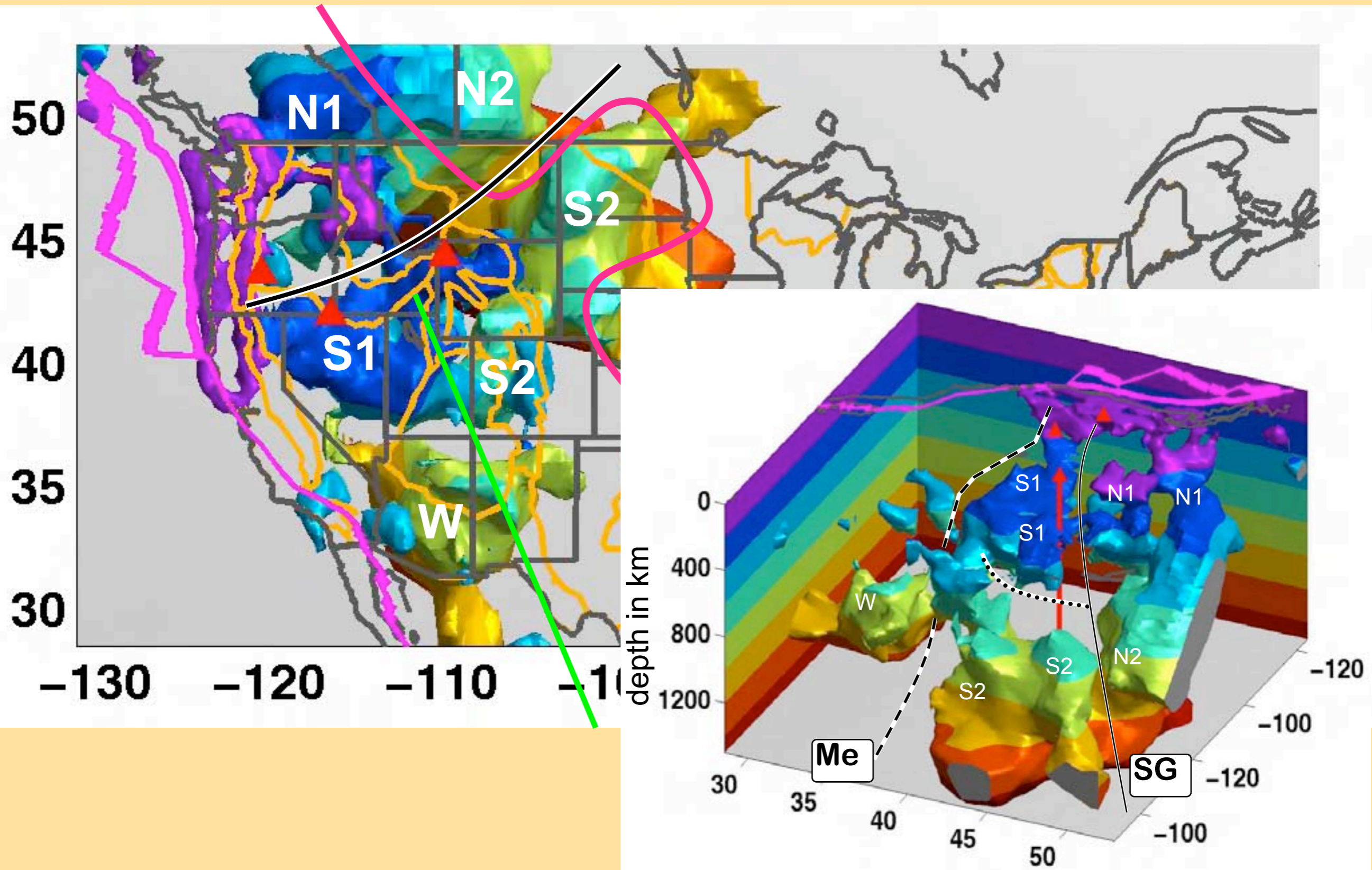


Sigloch et al., Nature Geosci., 2008

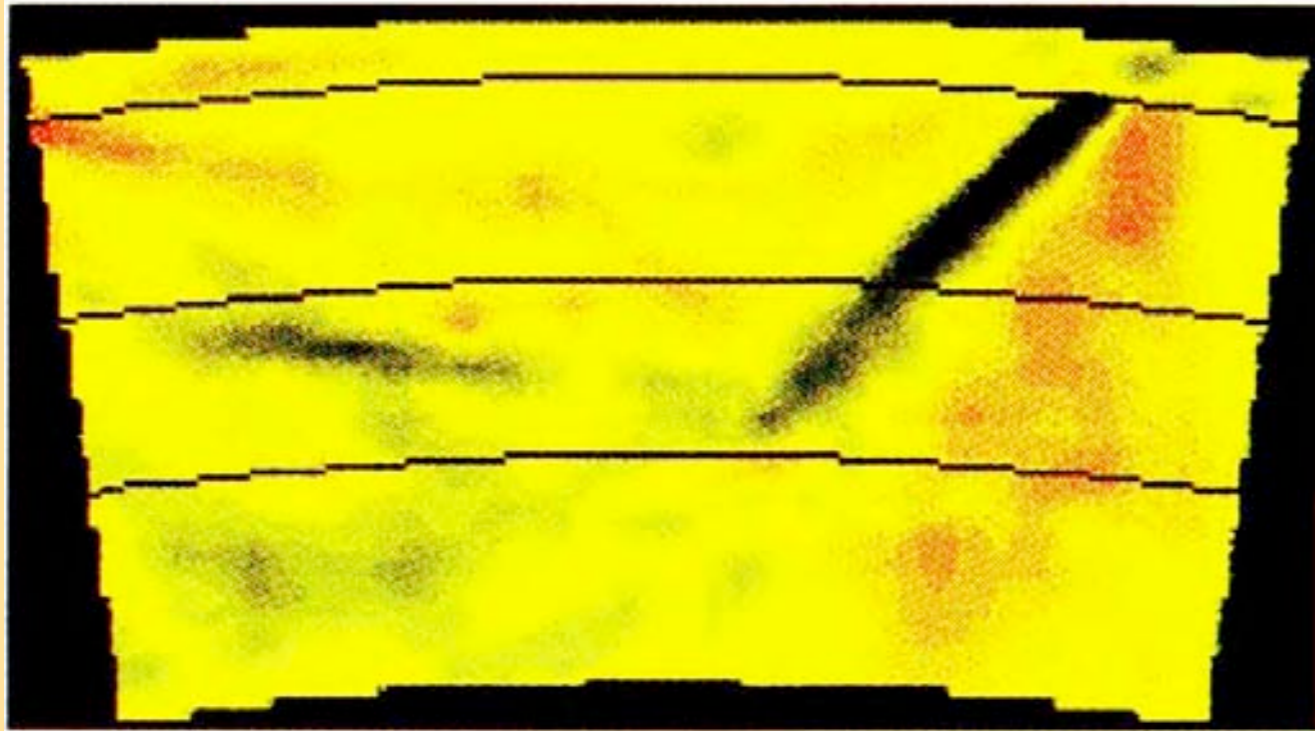
Slabs do not go gently into that good night



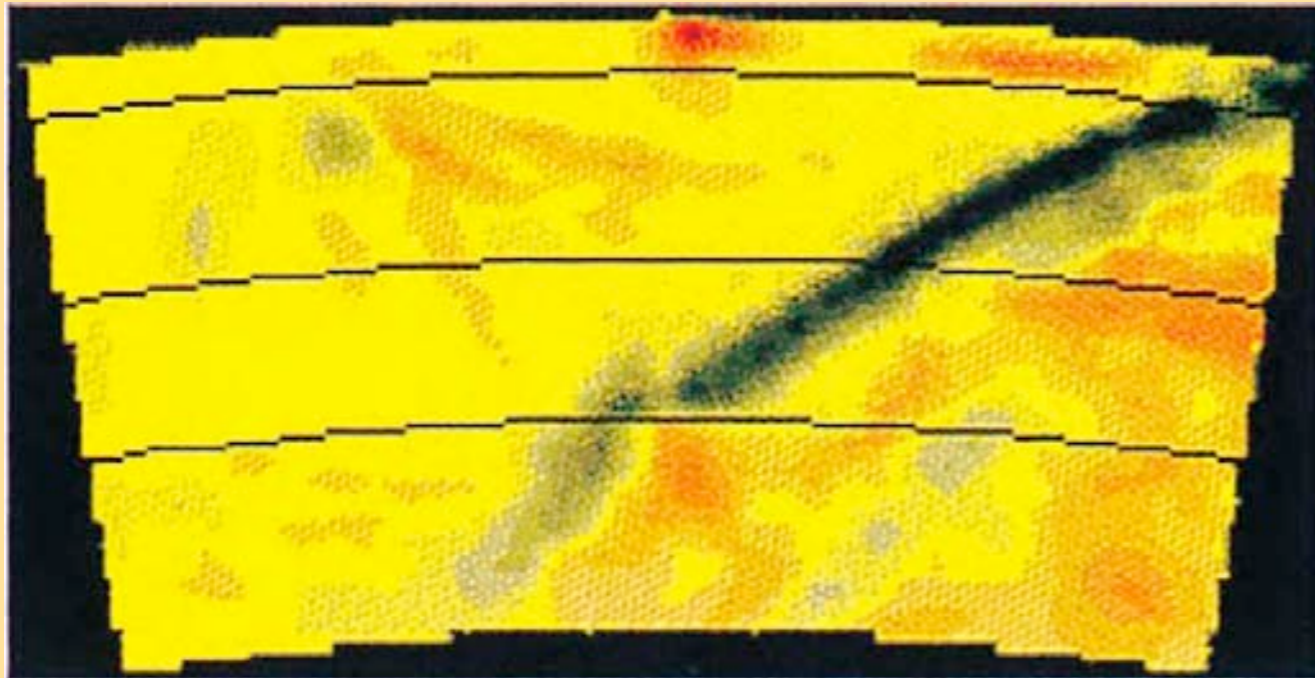
Slabs do not go gently into that good night



return flow across a TBL?

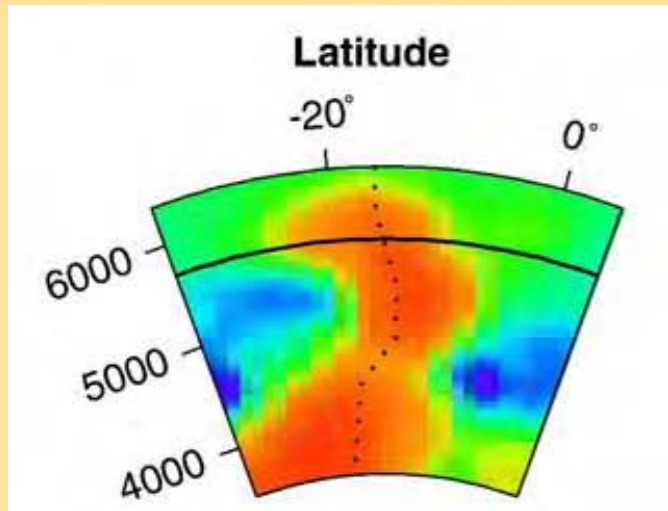


Tonga-Kermadec
(Deal et al., 1999)

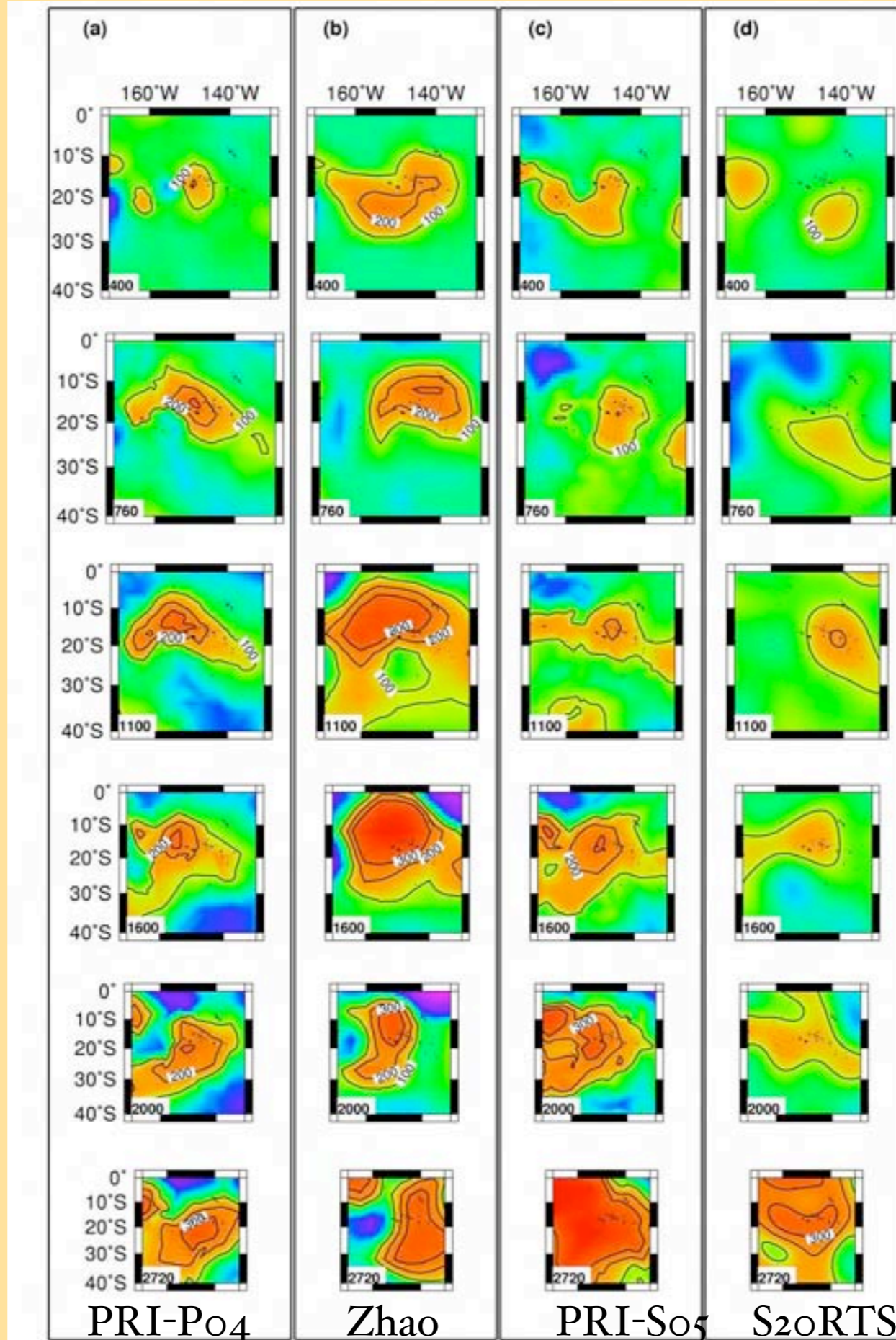


Japan 42° N
(Deal & Nolet, 1999)

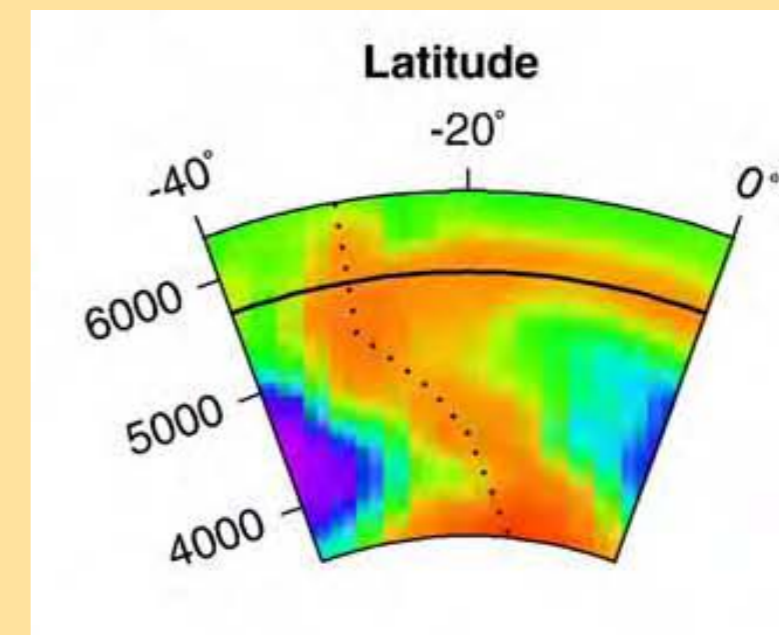
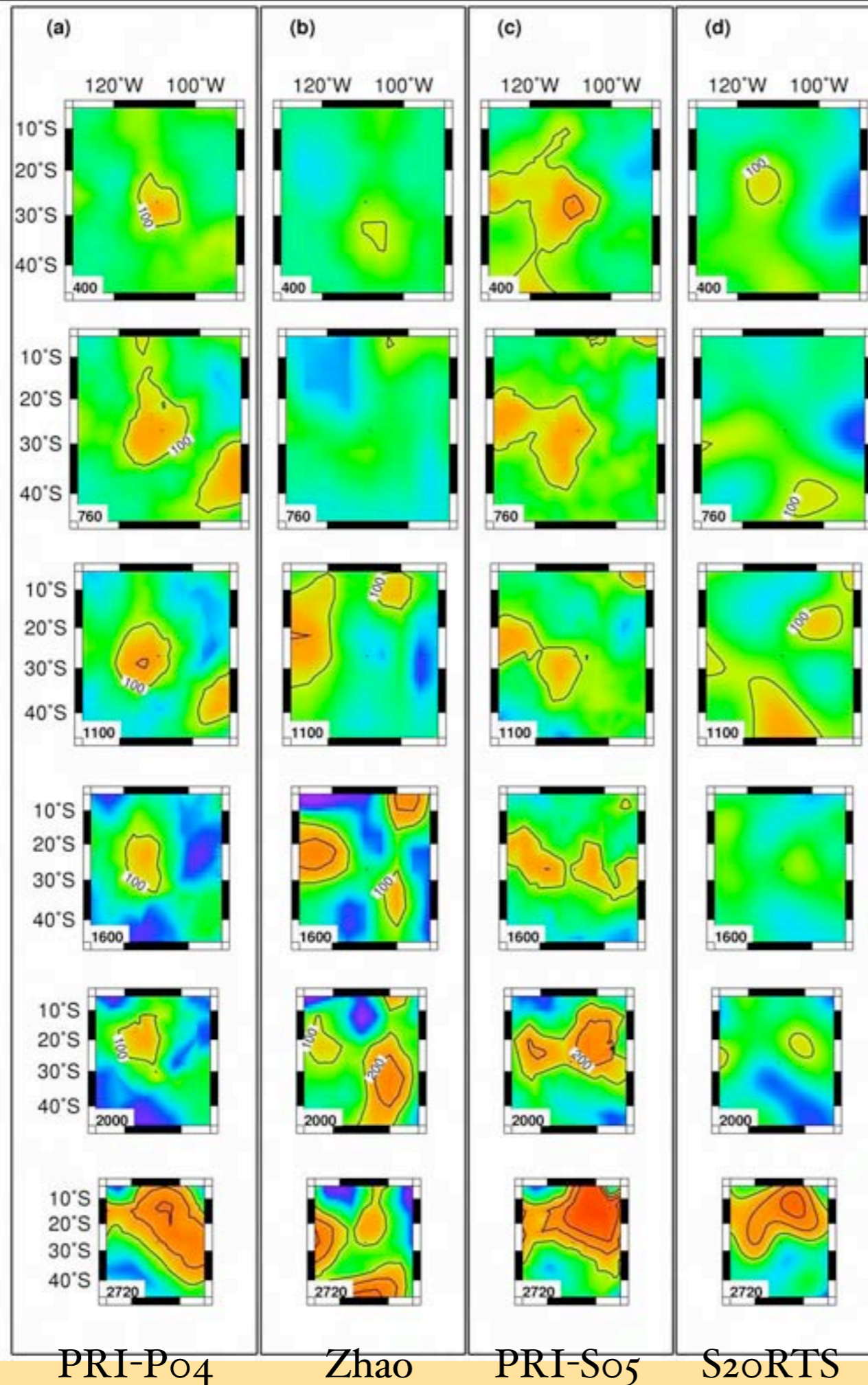




Tahiti
(Society Isl)



ΔV_P and ΔV_S
converted to
 ΔT (K)

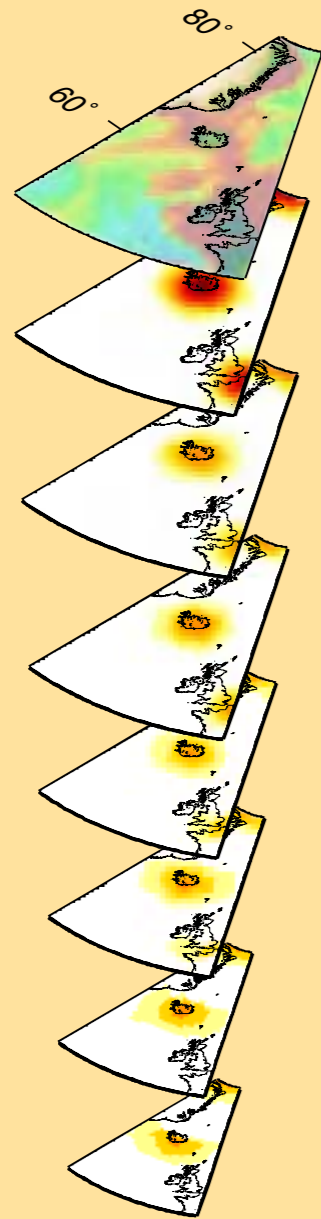


Easter island

ΔV_P and ΔV_S
 converted to
 ΔT (K)

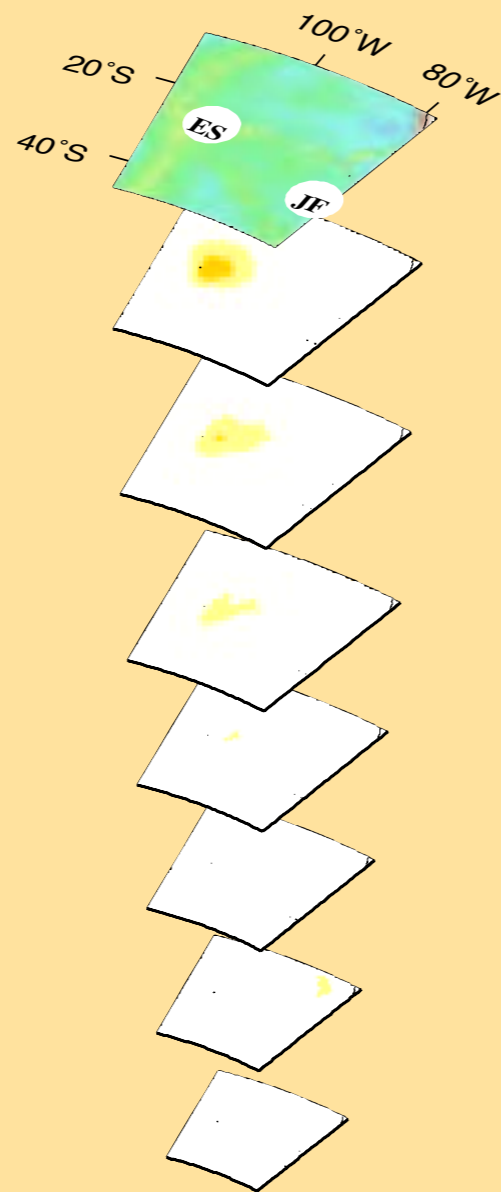
Thin plumes are not resolved

$z=2800$ km, $w=400$ km



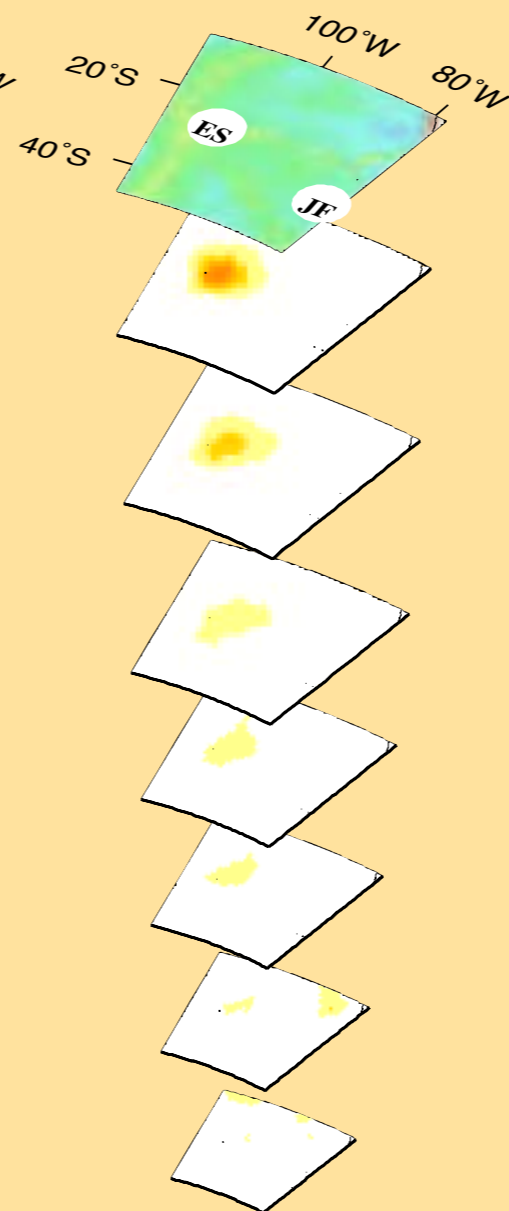
Input

Easter/Juan Fernandez



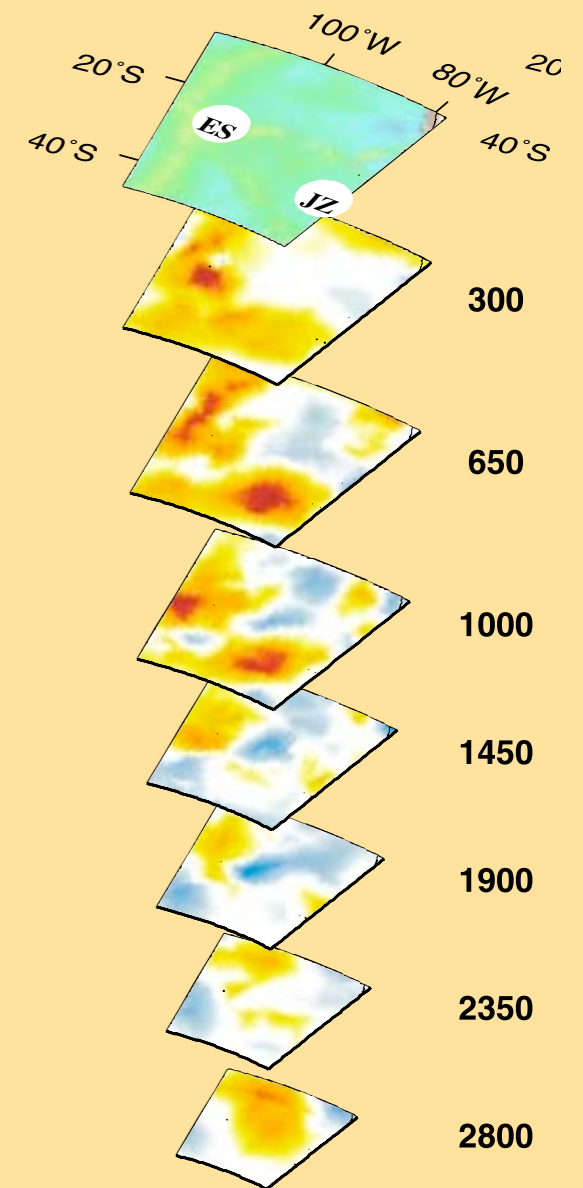
200 km

Easter/Juan Fernandez



400 km

P WAVES



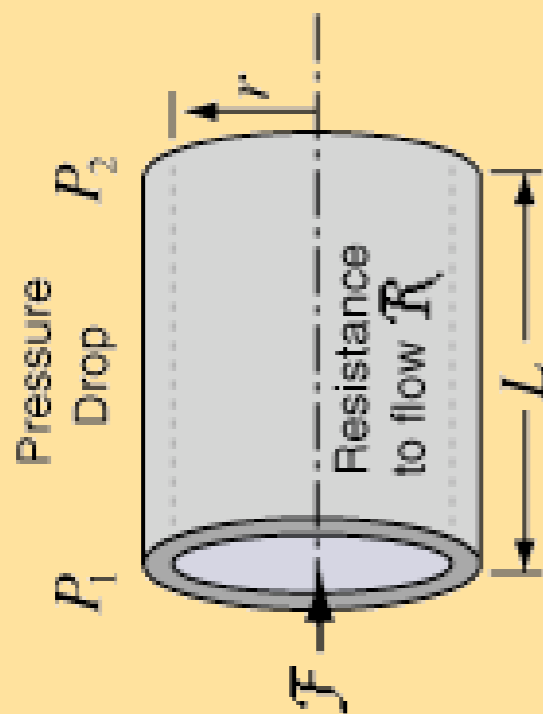
PRI-P05

Plumes are larger than we thought

- Buoyancy flux indicates weak plume flux
- Early estimates limit plume flux to ~ 3 TW
- Theory predicts narrow (<100 km) plumes
- But such plumes would *not* show up
- Plume width must be *several hundred km*

Can we estimate plume flux from tomography?

- The honest answer is: *barely*
- But even a simplified treatment leads to surprising lower limits in plume flux

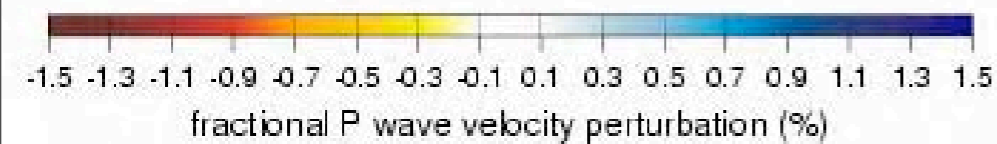
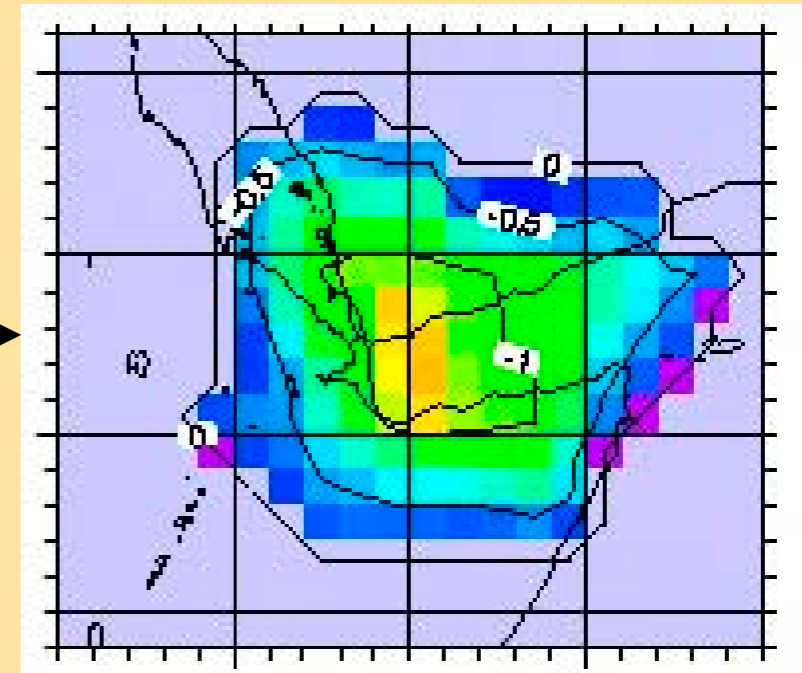
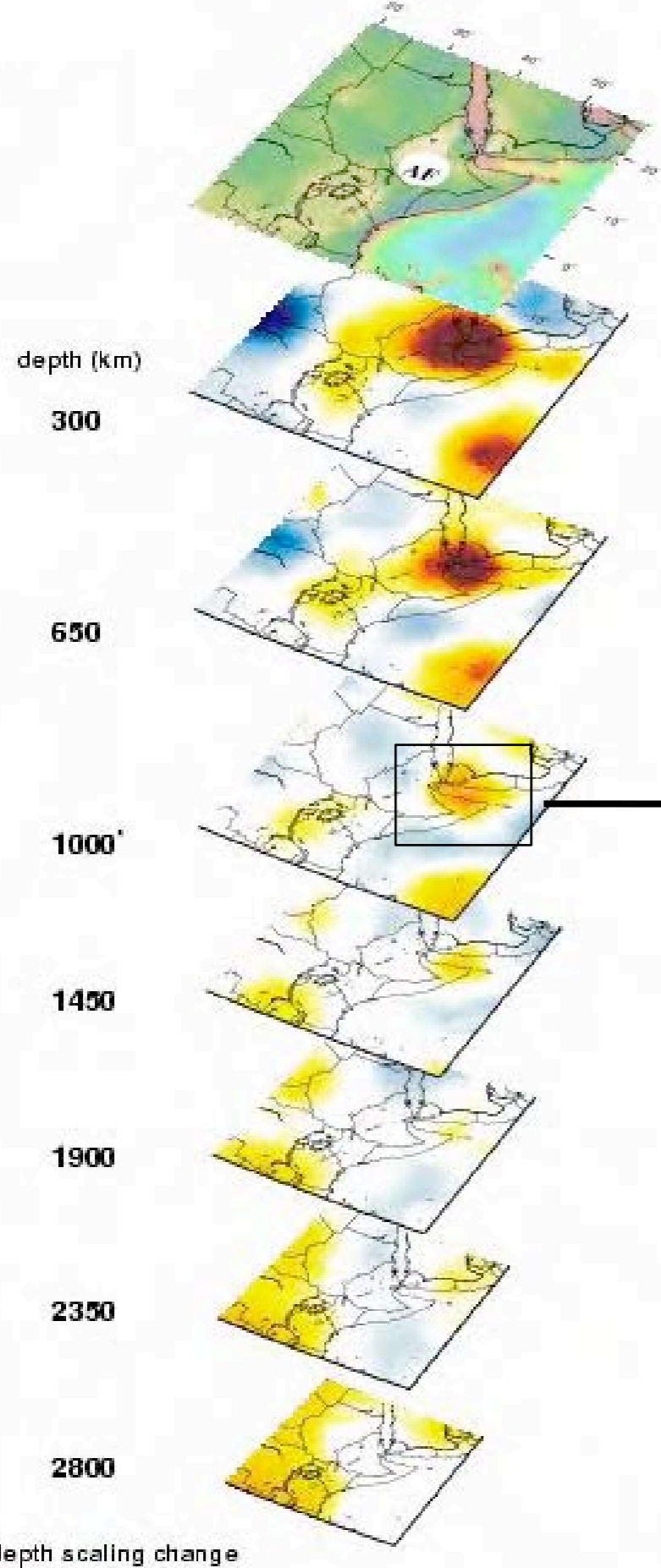


face of rotation whose meridian section is a parabola. In this case we speak of a *parabolic velocity profile*. The total volume W issuing per second is obtained by taking the integral $\int \mathbf{v} d\mathbf{S}$ over a cross-section. In this case we have

$$W = \int_0^a \frac{p_1 - p_0}{4\eta l} (a^2 - r^2) 2\pi r dr = \frac{\pi (p_1 - p_0) a^4}{8\eta l}. \quad (74)$$

This is Poiseuille's Formula, which states that the quantity of fluid issuing each second is directly proportional to the pressure difference and to the fourth power of the radius of the tube, and inversely pro-

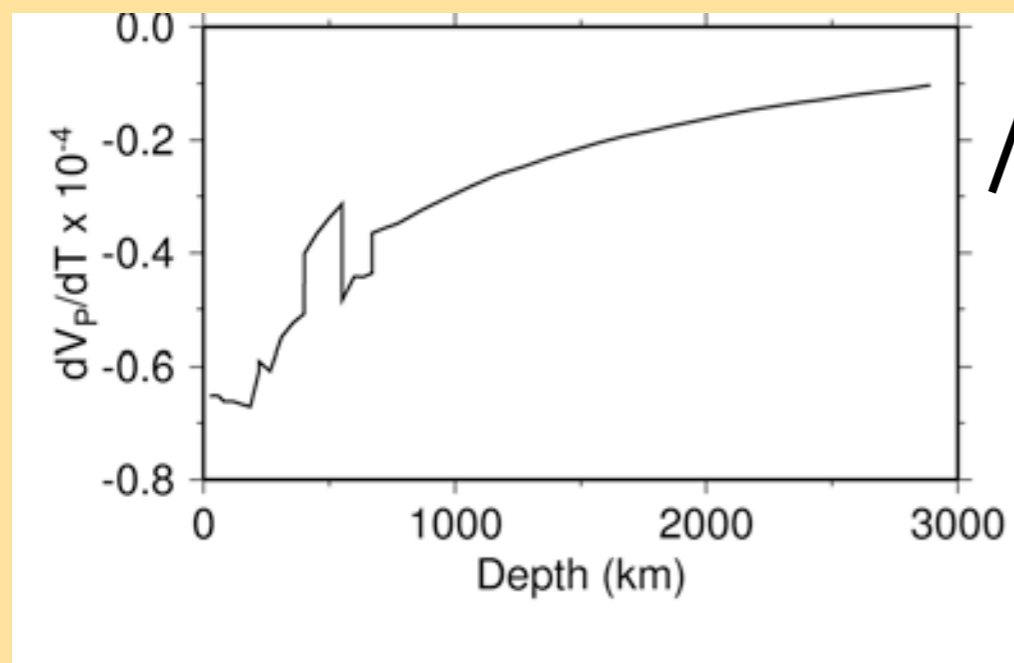
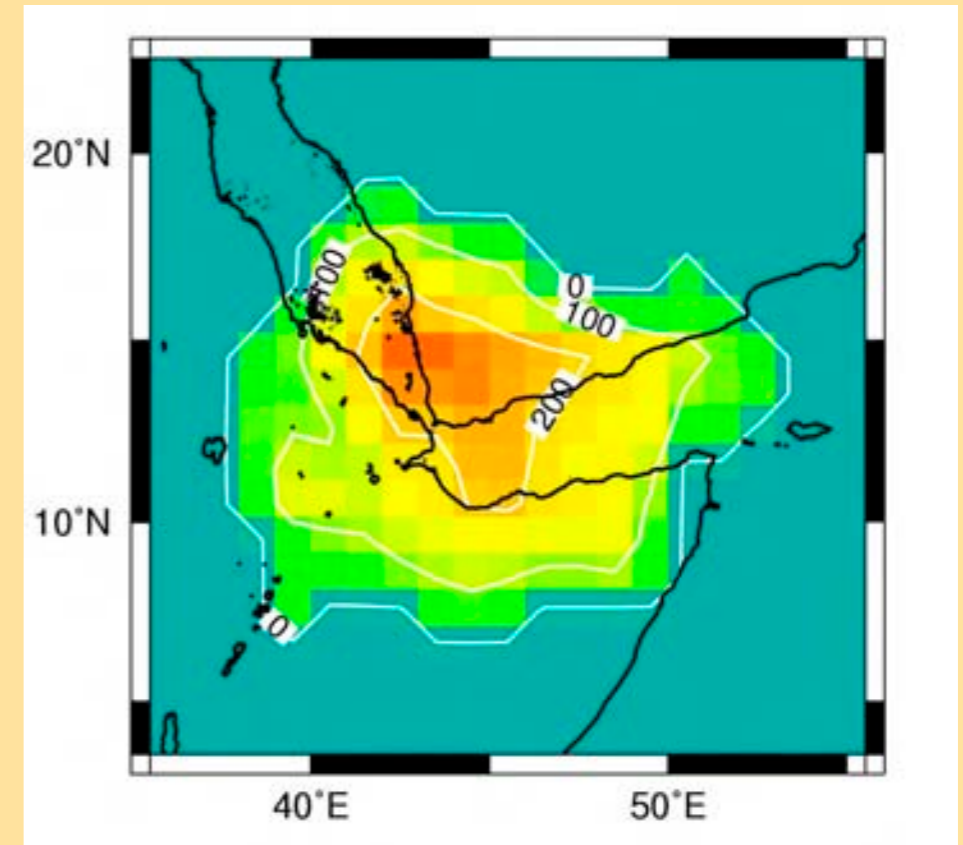
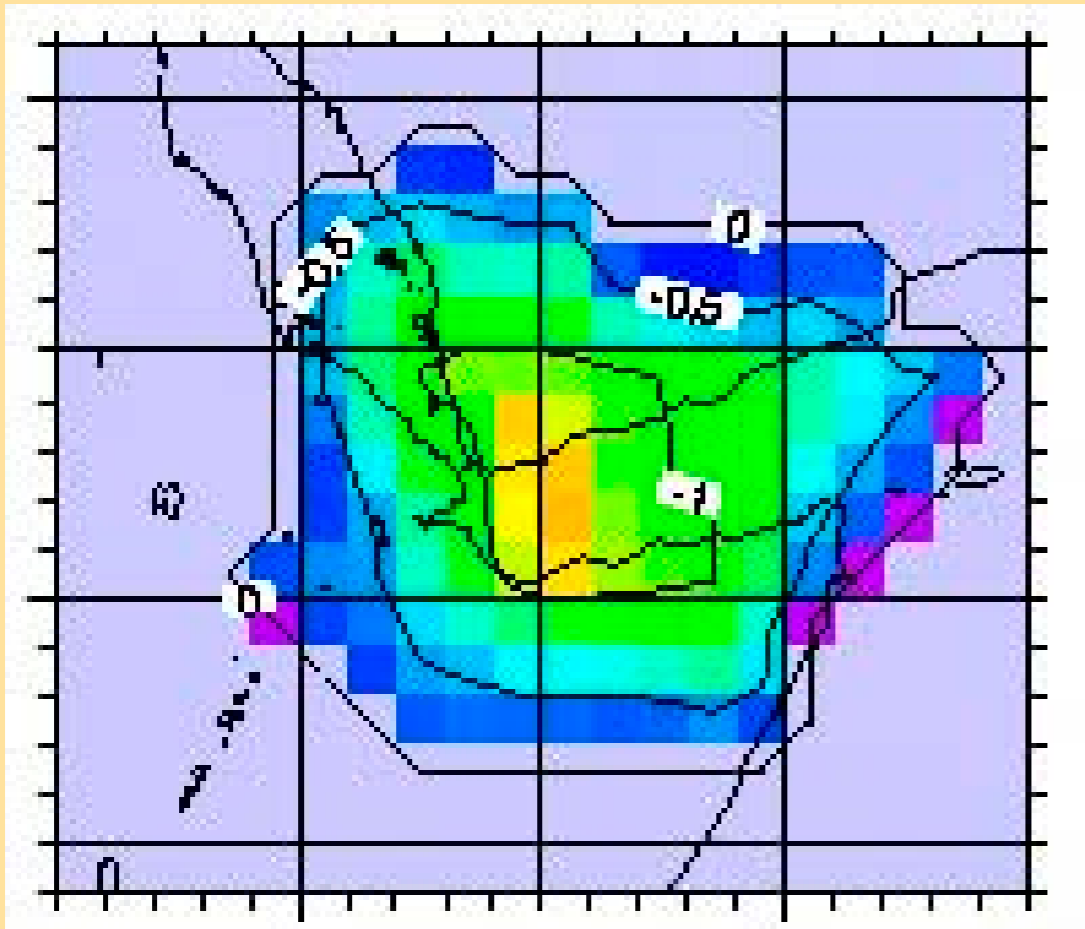
The method, step 1:
isolate a depth section from the
3D model of $\Delta V_p/V_p$



step 2: from V_P anomaly

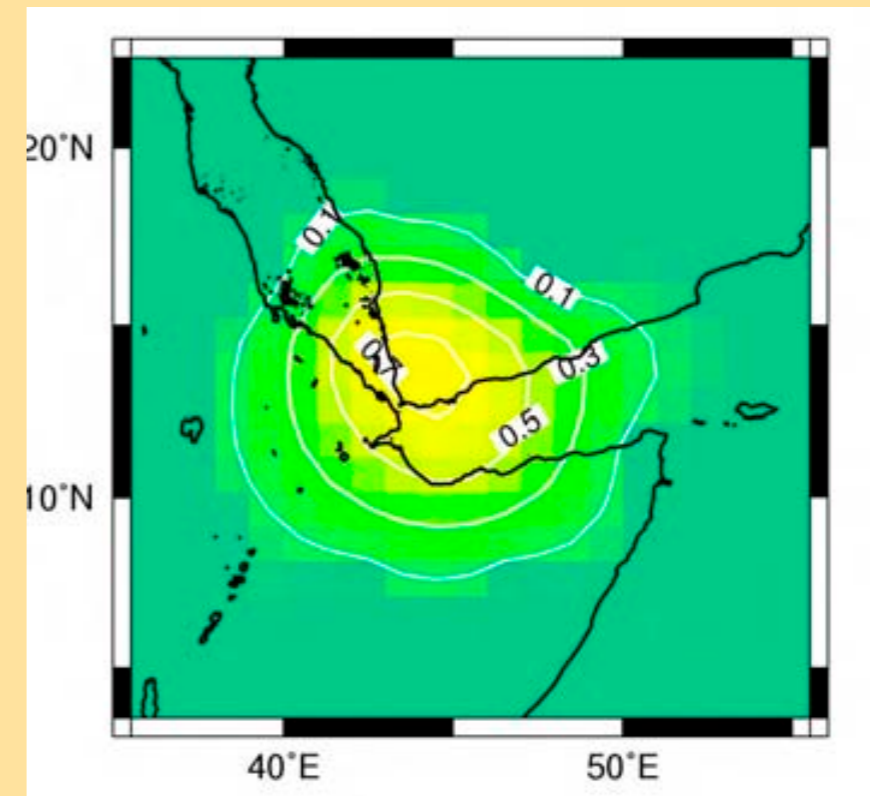
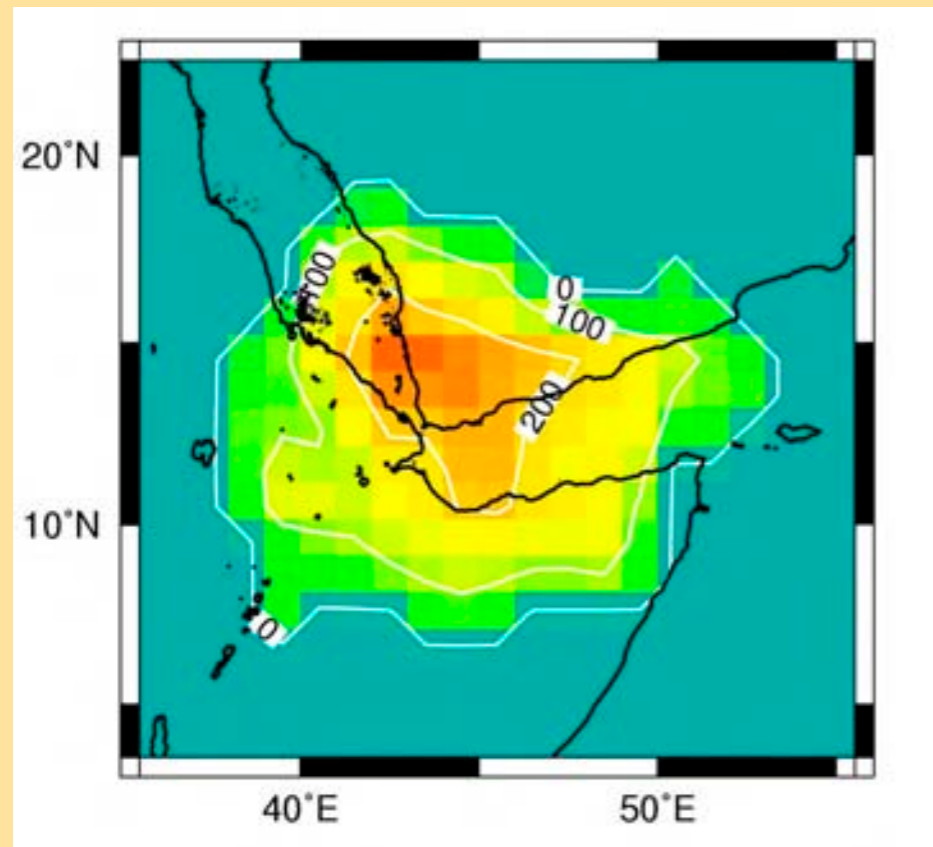
to

Temperature anomaly



3: from Temperature anomaly to

Rise velocity

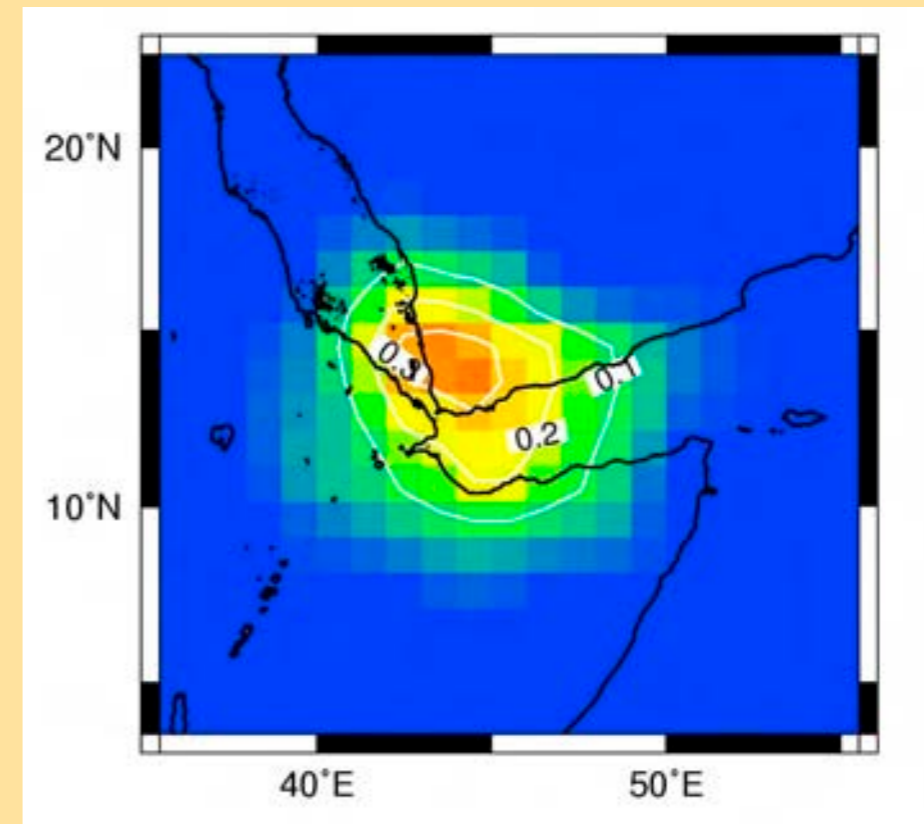
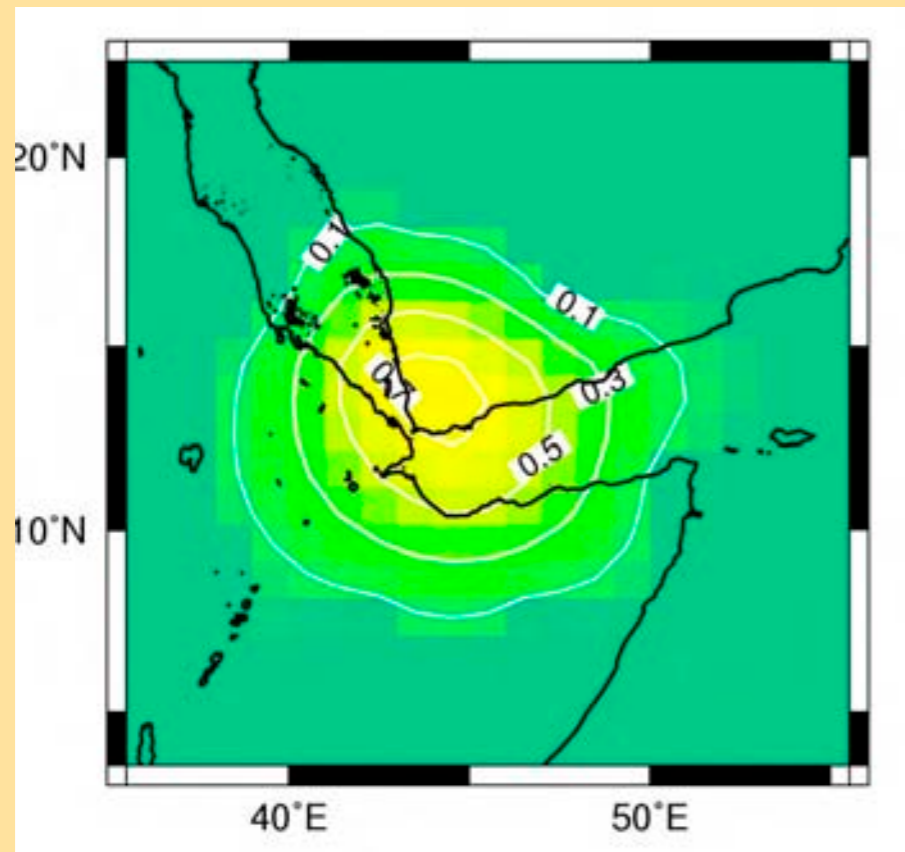


$$\eta \nabla^2 v_z = -\alpha \rho g \Delta T + g \Delta \rho F_e$$

$$\eta = \eta_R \exp\left(\frac{\beta T_m}{T}\right)$$

4: from rise velocity to

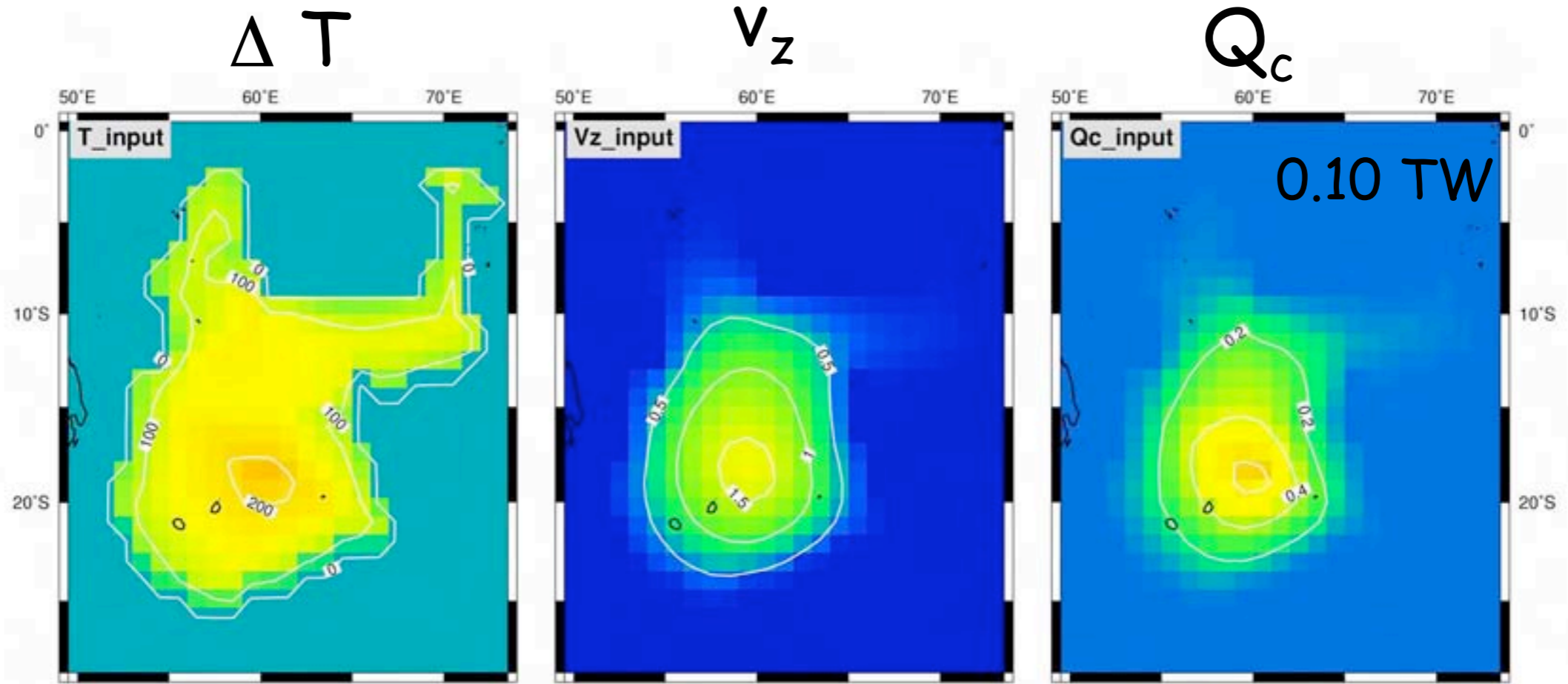
Heat flux



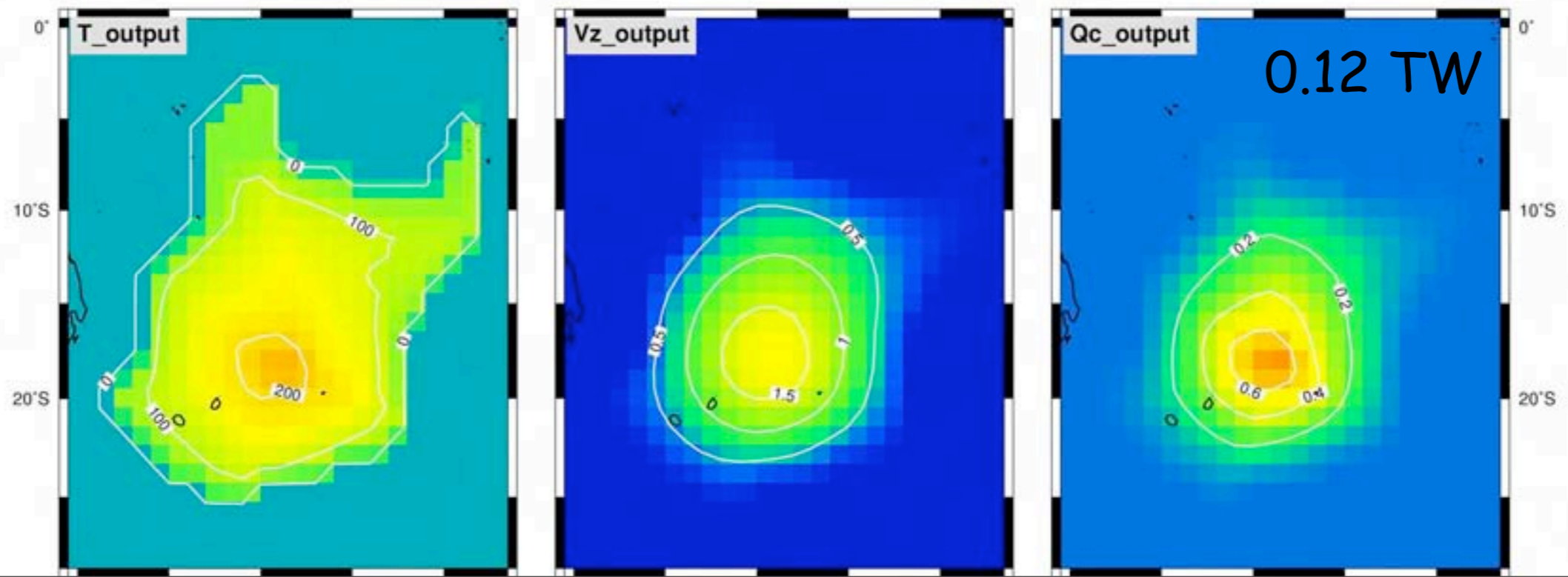
$$Q_c = c_p \rho \Delta T v_z$$

Sensitivity test: how good is the estimate?

in:

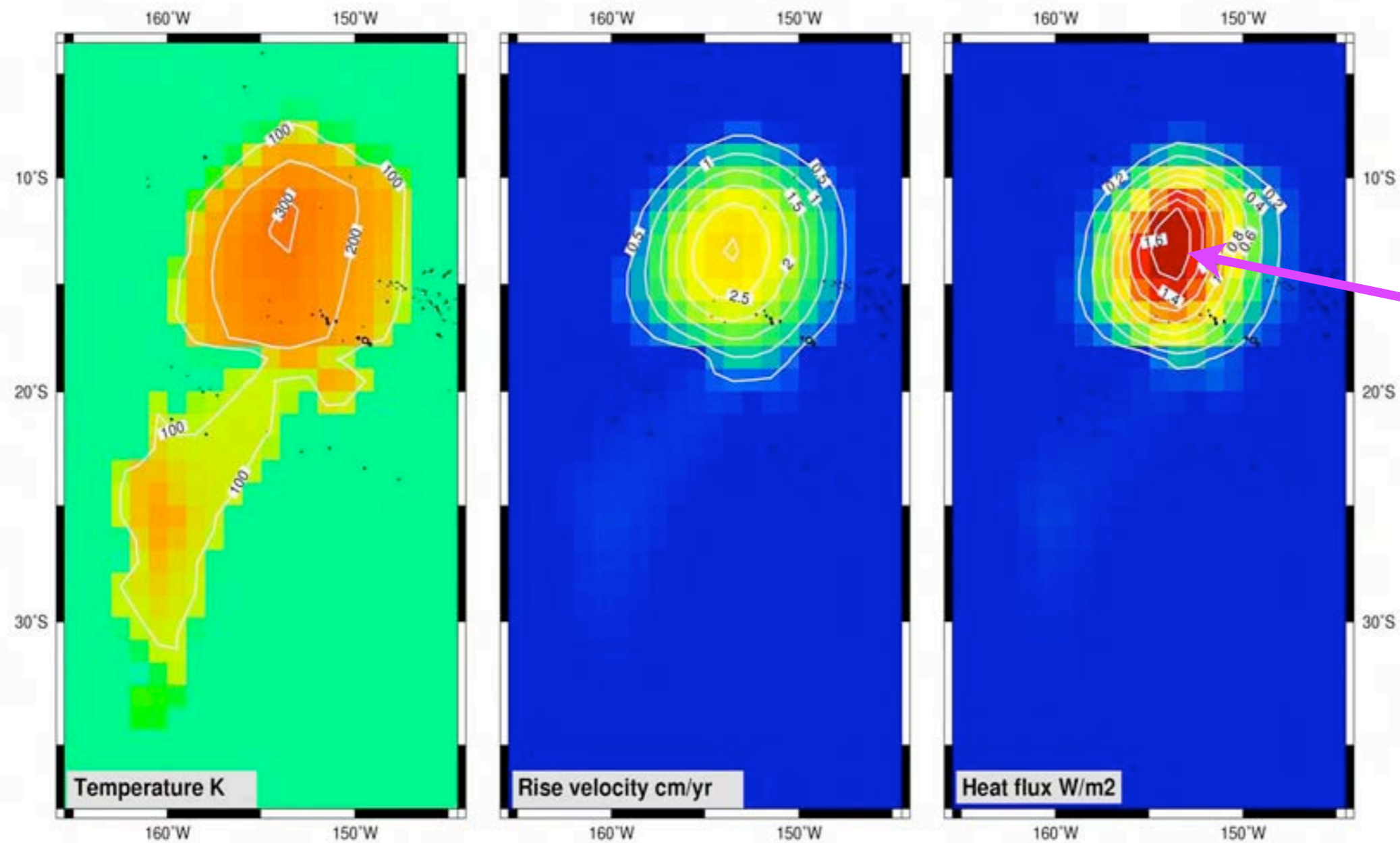


out:



Reunion
1100 km

Plumes are larger than we thought

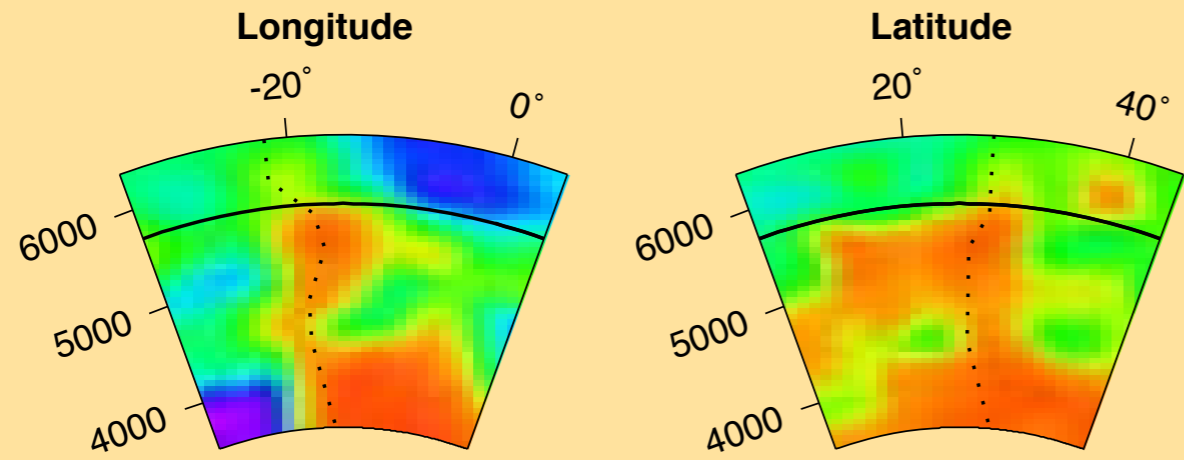


Total
flux
0.5 TW

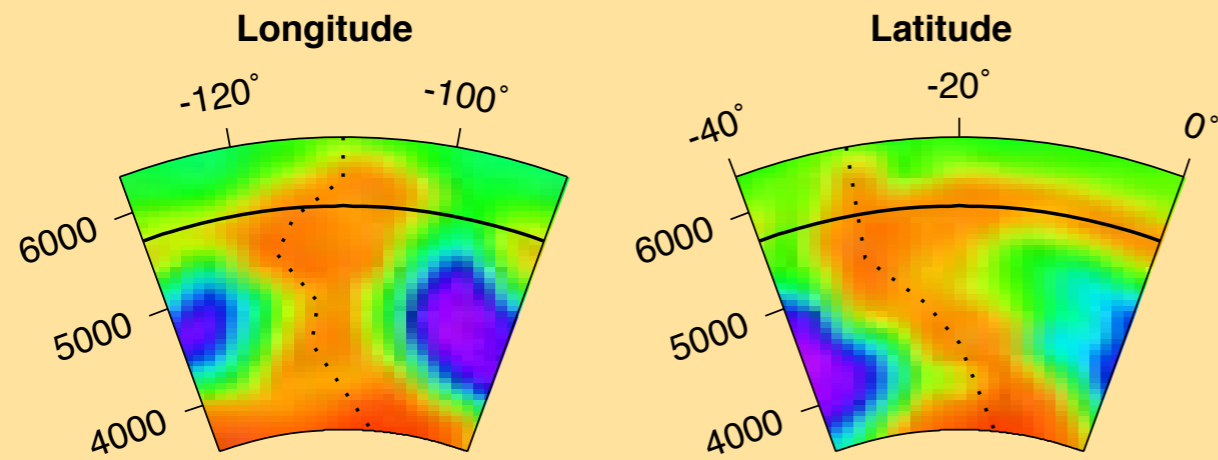
Tahiti plume at 1600 km depth (Nolet, EPSL 2006)

(Some) plumes
broaden below
660 km depth

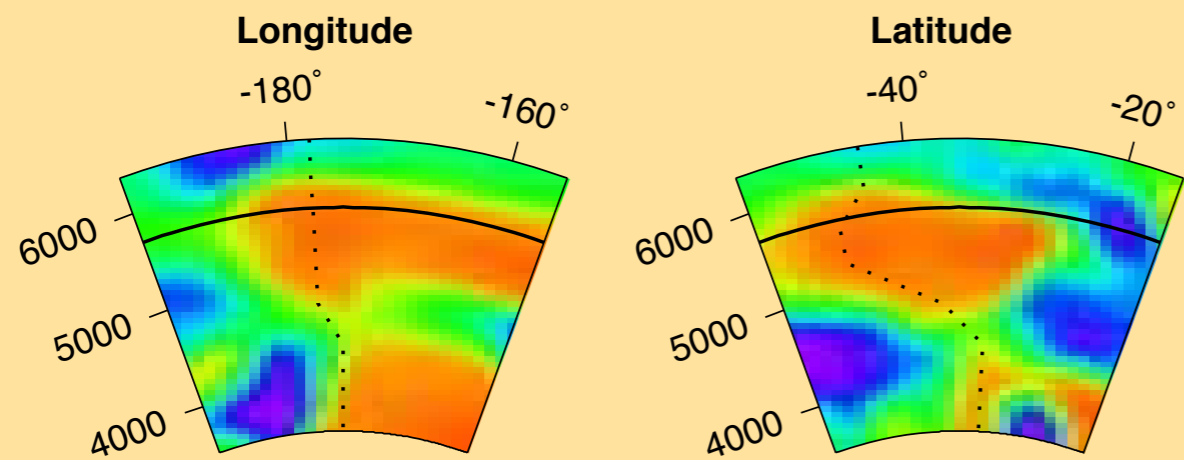
Canary



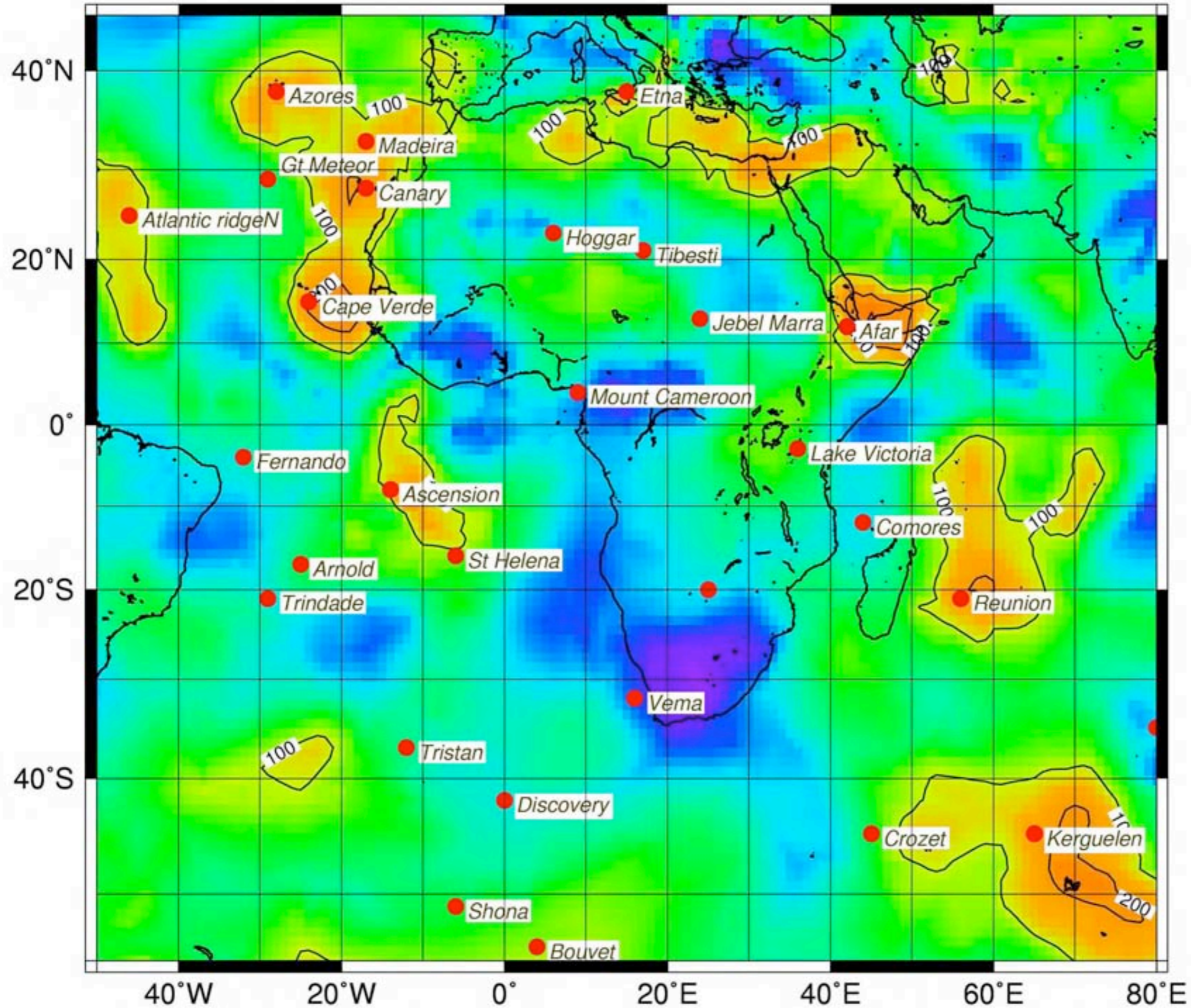
Easter



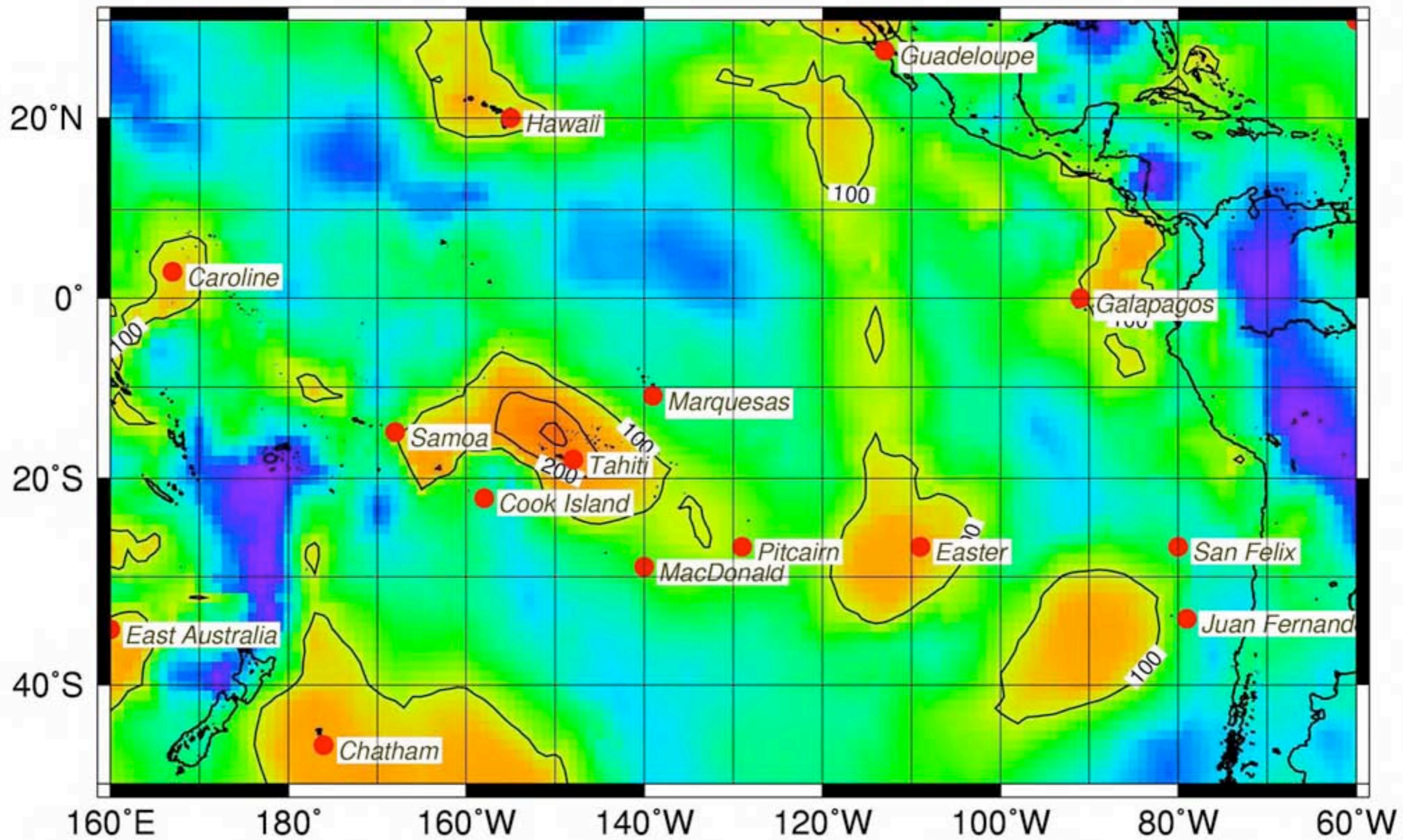
Chatham



Nolet et al., 2006



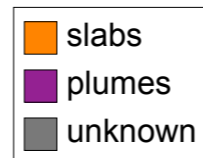
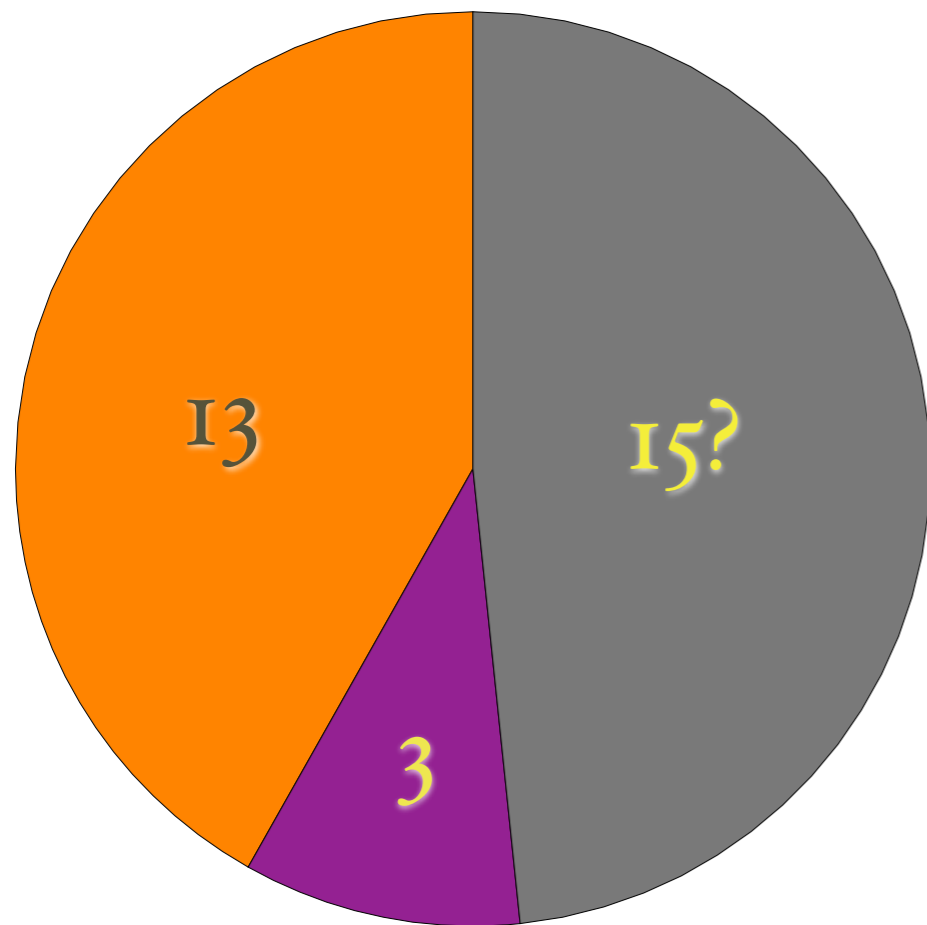
ΔT (K)
at 800 km



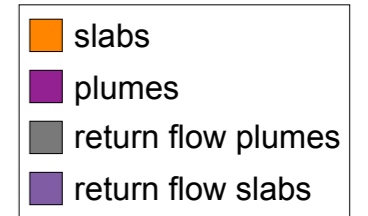
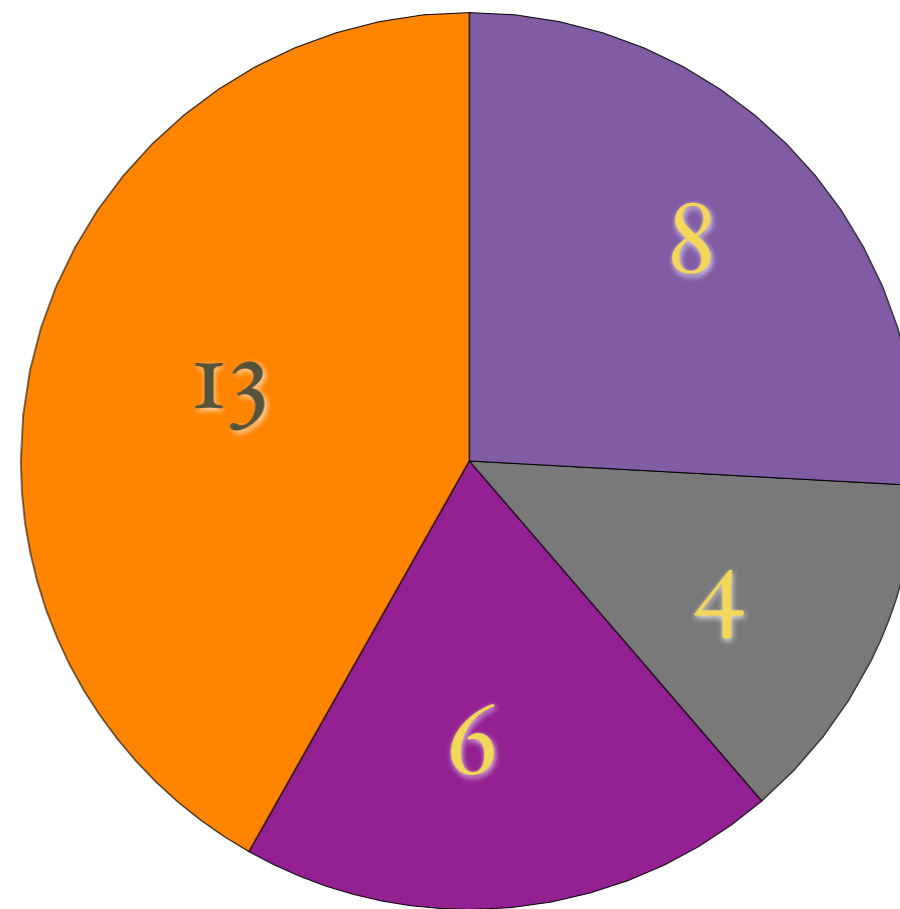
Inferred ΔT (K) from ΔV_p at 800 km depth

31 TW across 660 ?

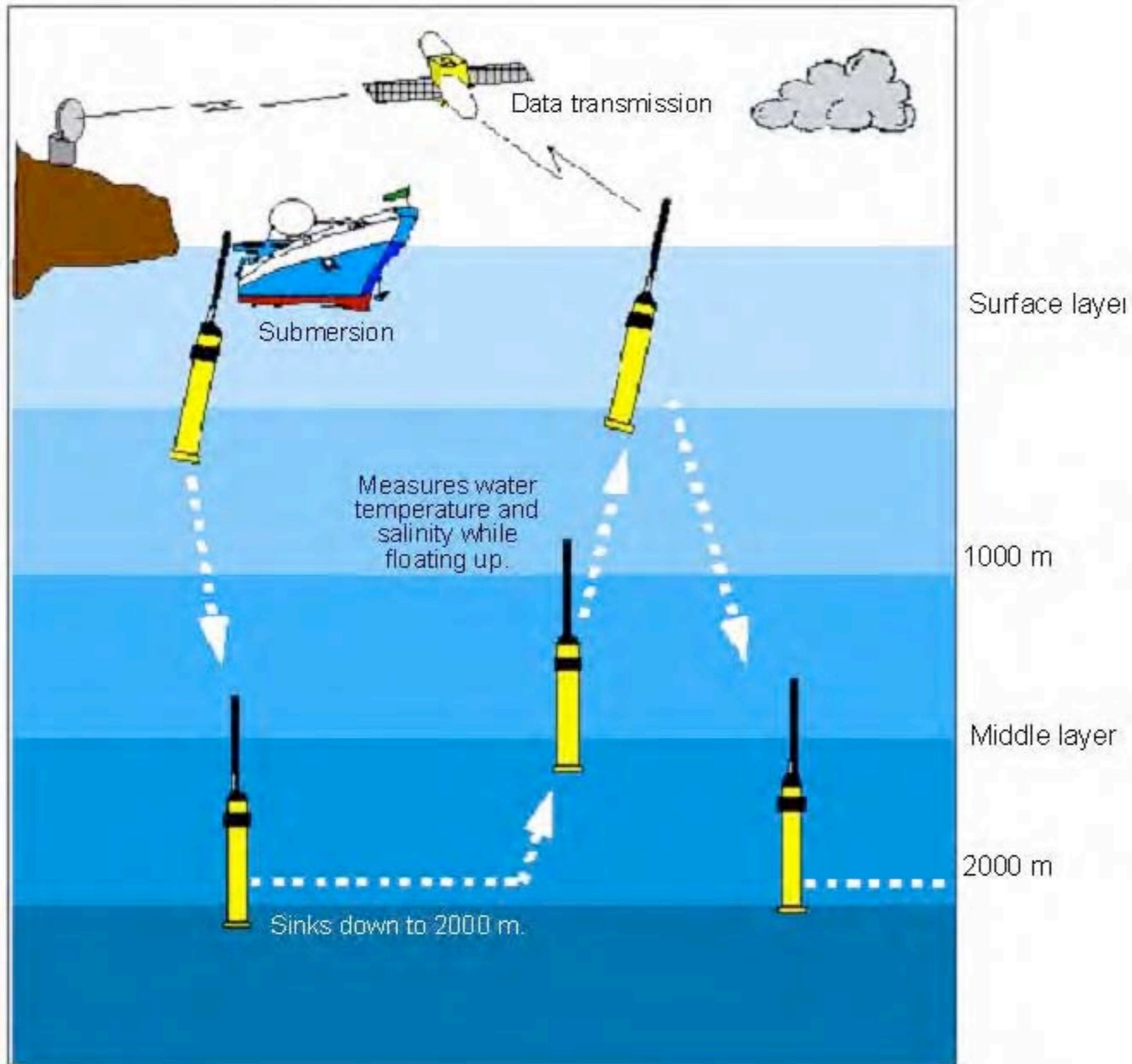
Classic view

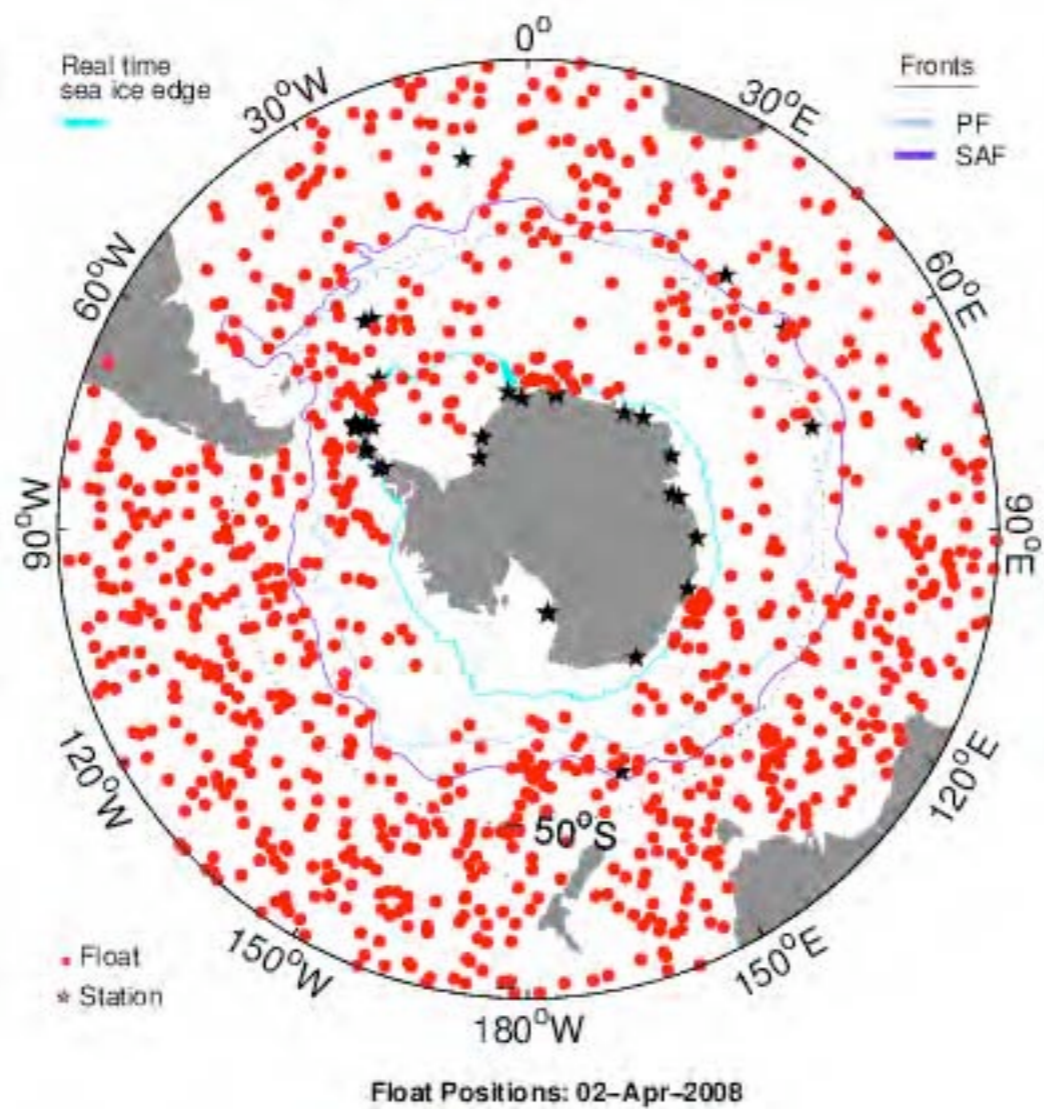


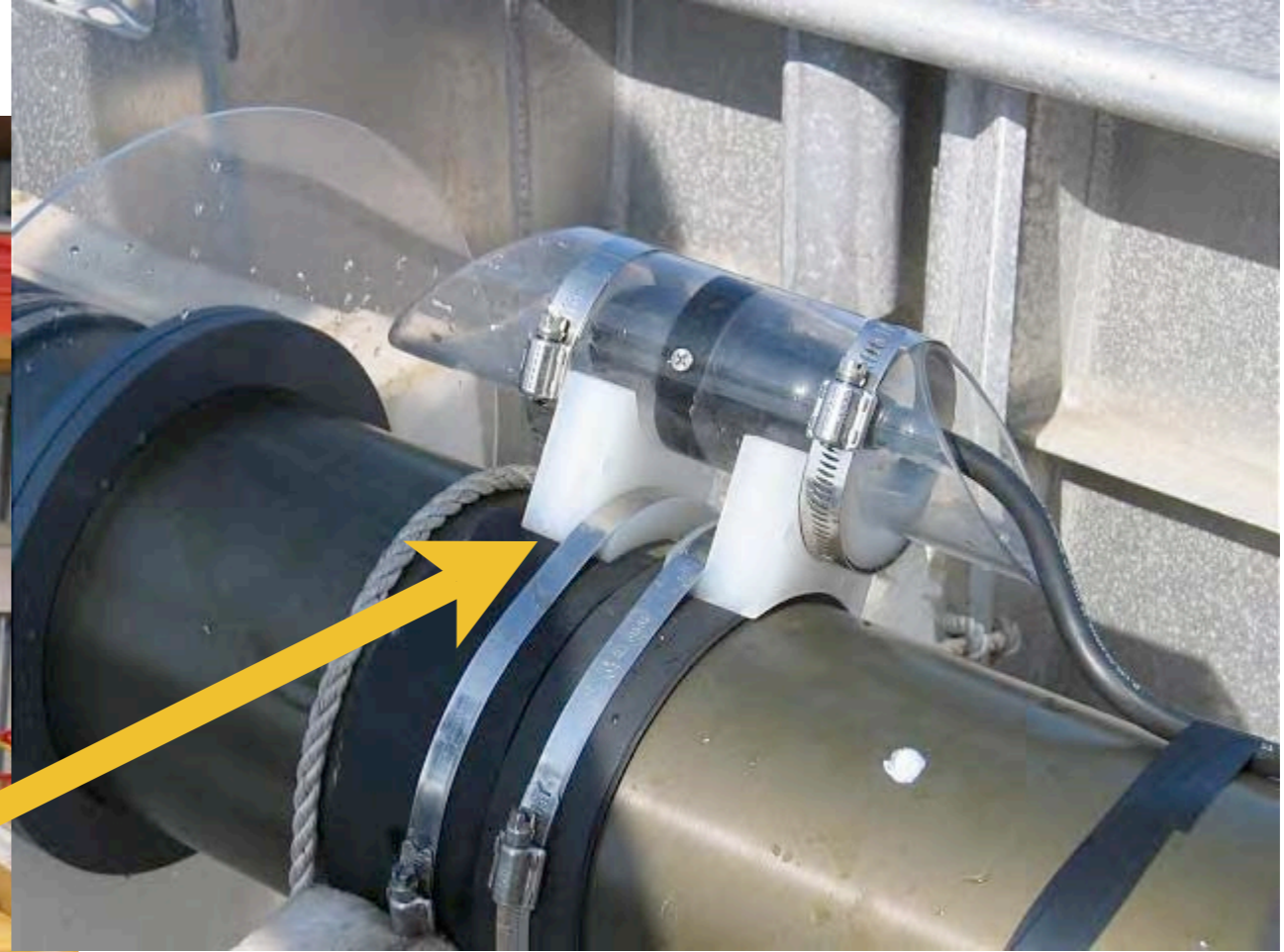
New view



ARGO floats

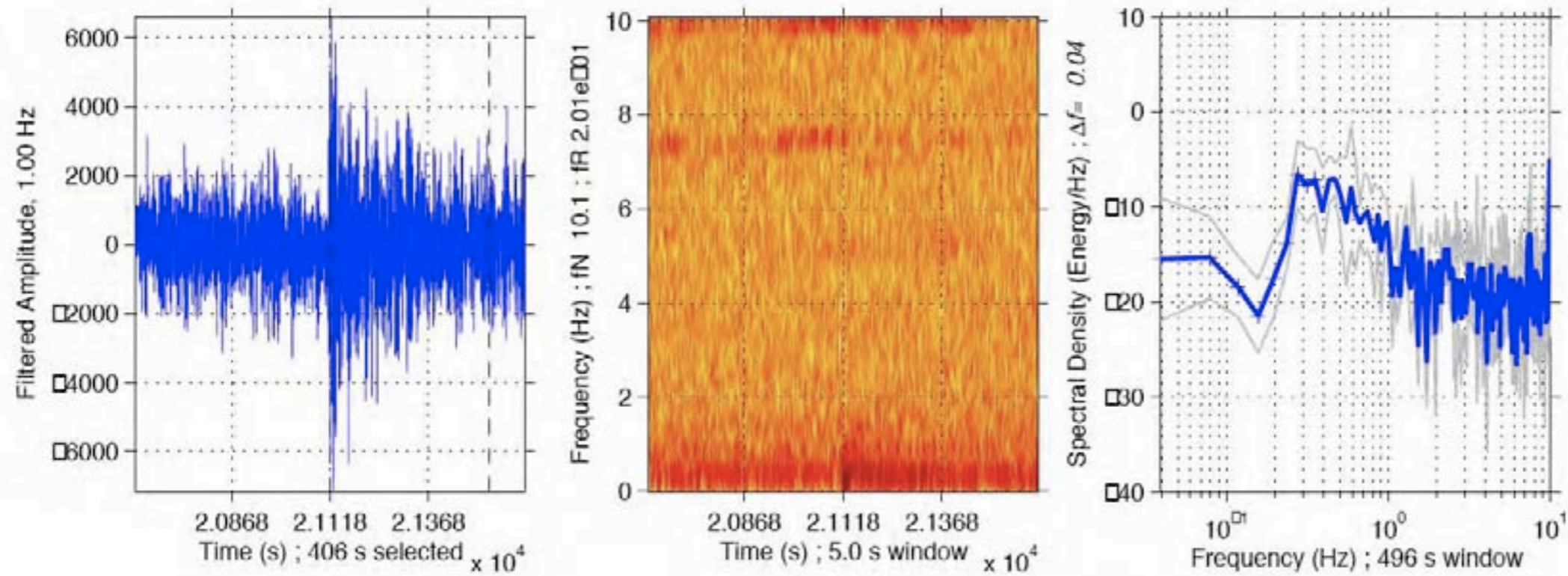






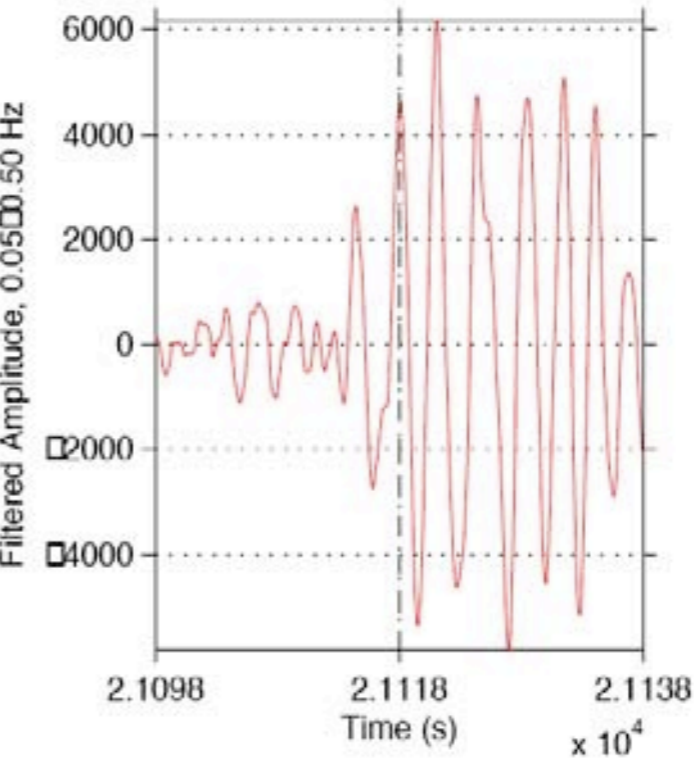
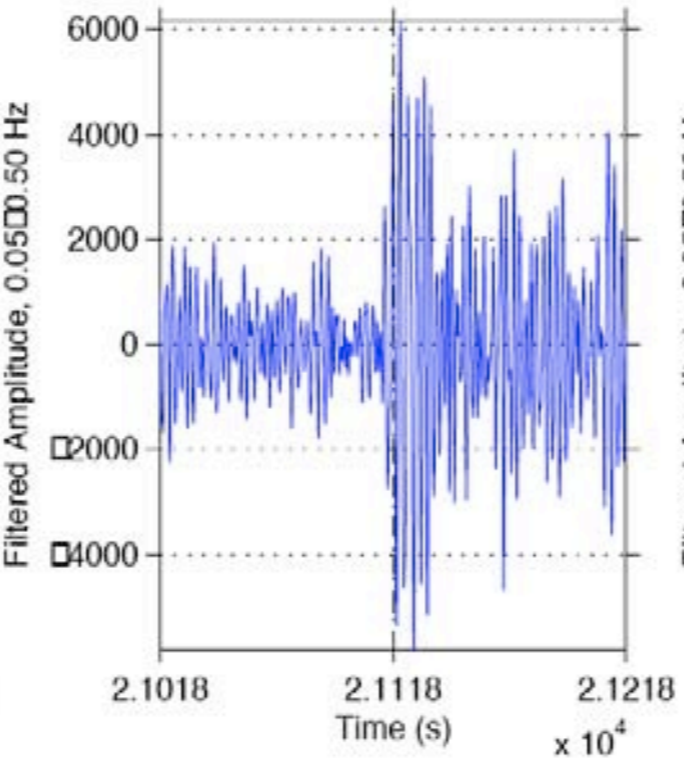
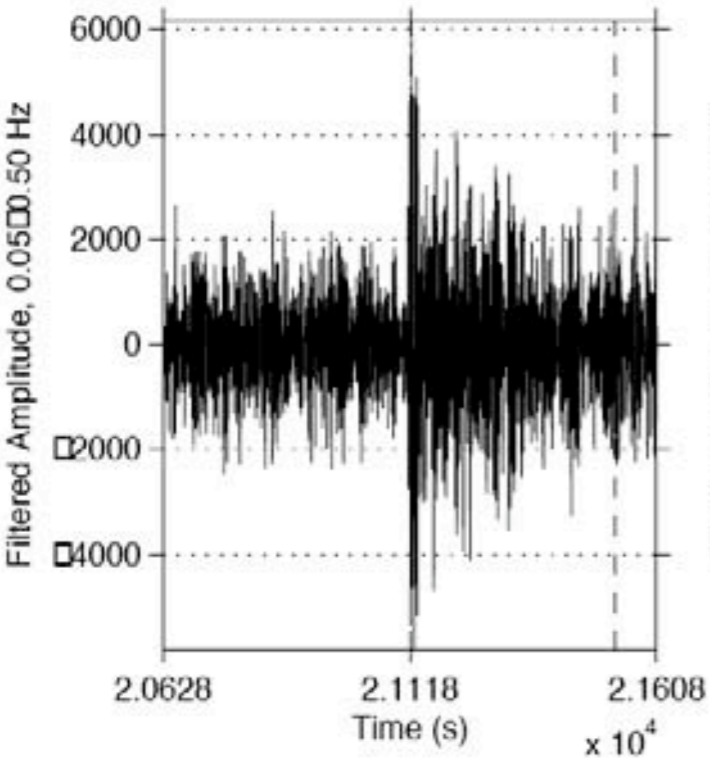
Frederik Simons (pers. comm.)

Near West Coast of Colombia Mw 5.95 at 46.5°



Frederik Simons (pers. comm.)

Zooming in on the onset \longrightarrow



Frederik Simons (pers. comm.)

Conclusion

- Slabs stalling at 660,
 - Plumes widening below 660,
 - Plume flux too large in lower mantle,
 - Low velocities below the slabs,
-
- TOMOGRAPHIC EVIDENCE POINTS TO THE 660 BEING A THERMAL BOUNDARY LAYER
 - MASS EXCHANGE: ONLY SLABS AND PLUMES